

CONDENSED MATTER PHYSICS

Gate-tuned quantum Hall states in Dirac semimetal $(\text{Cd}_{1-x}\text{Zn}_x)_3\text{As}_2$ Shinichi Nishihaya,¹ Masaki Uchida,^{1*} Yusuke Nakazawa,¹ Markus Kriener,² Yusuke Kozuka,¹ Yasujiro Taguchi,² Masashi Kawasaki^{1,2}

The recent discovery of topological Dirac semimetals (DSMs) has provoked intense curiosity not only regarding Weyl physics in solids but also about topological phase transitions originating from DSMs. One specific area of interest is controlling the dimensionality to realize two-dimensional quantum phases such as quantum Hall and quantum spin Hall states. For investigating these phases, the Fermi level is a key controlling parameter. From this perspective, we report the carrier density control of quantum Hall states realized in thin films of DSM Cd_3As_2 . Chemical doping of Zn combined with electrostatic gating has enabled us to tune the carrier density both over a wide range and continuously, even across the charge neutrality point. Comprehensive analyses of gate-tuned quantum transport have revealed Landau-level formation from linearly dispersed sub-bands and its contribution to the quantum Hall states. Our findings also pave the way for investigating the low-energy physics near the Dirac points of DSMs.

INTRODUCTION

In the past decade, electron transport in Dirac systems has attracted intensive research efforts to uncover various kinds of unprecedented phenomena in two-dimensional (2D) Dirac systems such as graphene (1) and the surface states of topological insulators (2–4). Recently, this field has grown to incorporate even 3D systems after the experimental discovery of topological Dirac semimetals (DSMs) (5–7) and Weyl semimetals (8–10) in solids, where the bulk conduction band (CB) and valence band (VB) touch at finite pairs of Dirac or Weyl points to form 3D Dirac dispersions. Among the DSM materials, Cd_3As_2 has been treated as an ideal test bed because its ultrahigh mobility and chemical stability in air make it rich in potential for finding new topological phases. Soon after its theoretical prediction (5), the characteristic linear dispersion and the band-inverted nature of DSMs have been revealed by spectroscopy experiments (6, 11, 12). They have been subsequently followed by a surge of reports on transport phenomena including quantum oscillations (13–15), chiral anomaly (16–18), and Fermi arc-mediated surface transport (19).

Although the above experimental observations on the transport of DSMs are mainly based on the 3D nature of bulk samples, a 2D confined Fermi surface achievable in a thin-film sample also offers opportunities for realizing novel 2D quantum states. The recent observation of the quantum Hall effect (QHE) in Cd_3As_2 is one example (20–22). It has become apparent that, due to the large sub-band splitting originating from the linear dispersion of Cd_3As_2 , a 2D nature of the Fermi surface can easily emerge even with a moderate confinement (20). This aspect of Cd_3As_2 is highly beneficial for realizing another type of 2D quantum state called quantum spin Hall insulator, which we predict to be achievable within the band-inverted regime of Cd_3As_2 (5, 20, 23).

From the viewpoint of transport measurements, however, the major obstacle in the pursuit of quantum phases in Cd_3As_2 , whether using 3D bulk or 2D confined films, is the defect-induced high carrier density, which shifts the Fermi level away from the charge-neutral Dirac points. Because Cd_3As_2 is easily electron-doped due to As deficiencies incorporated during the synthesis processes, achieving a low electron

density still remains challenging even in samples with high crystallinity (13–15, 19, 20). Because the band properties accessible through transport measurements are limited to those at the Fermi level, systematic carrier control of Cd_3As_2 is of crucial importance to elucidating its electronic structure in detail. In this regard, chemical doping of Zn is known to be an effective way to reduce the electron density of Cd_3As_2 . By making a solid solution with Zn_3As_2 , a p-type trivial semiconductor with a band gap of 0.93 to 1.1 eV (24), the residual electrons of Cd_3As_2 can be effectively compensated (25, 26). The sole yet serious drawback of the use of Zn doping, though, is that it alters the band structure because heavy doping would lead to a suppression of the band inversion in Cd_3As_2 and hence loss of its DSM nature (26). In terms of avoiding this heavy doping effect, electrostatic carrier depletion in a field effect transistor configuration is also a useful choice in the case of thin-film samples. So far, Liu *et al.* (27) have reported gate-tuning using an electric double-layer transistor on rather thick Cd_3As_2 films (~50 nm). The capability of tuning the Fermi-level position would advance the study of 2D quantum transport and clarify the actual band structure of confined films as well.

In this context, we report quantum transport of 2D confined Cd_3As_2 films, where the carrier density is controlled over a wide range using a combination of Zn doping and electrostatic gating (Fig. 1A). The effective carrier suppression with minimal Zn doping makes it possible to keep the high mobility and linear dispersion of Cd_3As_2 . The 2D nature of the confined Fermi surface emerges as the electron density is reduced, leading to the observation of clear quantum Hall (QH) states with various filling factors. Detailed analysis of the QH states while tuning the Fermi-level position has revealed Landau-level formation from linearly dispersed sub-bands and their occupation transitions depending on the Fermi-level position, which, in turn, provides a clearer understanding of the band structure and the confinement effect in Cd_3As_2 films.

RESULTS

Zn doping and field effect in Cd_3As_2 thin films

We fabricated Zn-doped Cd_3As_2 thin films on SrTiO_3 (100) substrate by a combination of pulsed laser deposition and thermal treatment, following the same procedures explained by Uchida *et al.* (20) (see also Materials and Methods). From x-ray diffraction measurements, the

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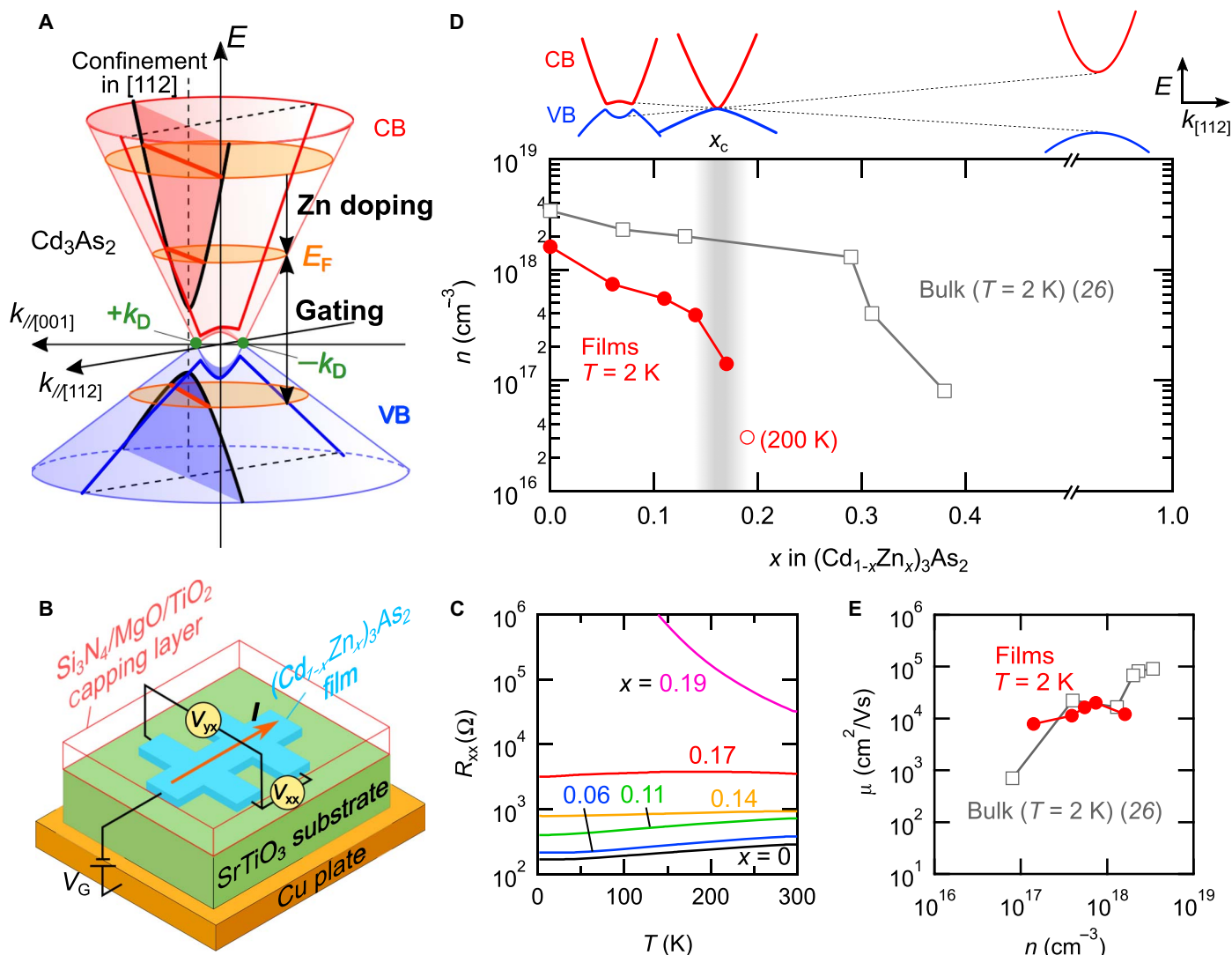


Fig. 1. Carrier control in Cd_3As_2 thin films by Zn doping and gating. (A) Schematics of the band structure of 2D confined Cd_3As_2 . The bulk Dirac points ($\pm k_D$) are on the [001] axis, whereas the growth direction of the thin films is along the [112] axis. Quantum confinement induces the formation of quantized sub-bands on the plane perpendicular to the [112] axis, and the sub-band dispersion is described by a hyperbolic function. The carrier density is controlled by the combination of chemical Zn doping and electrostatic gating in a transistor configuration in order to investigate quantum transport at various Fermi-level positions (E_F). (B) Schematic of back-gate transistor configuration using SrTiO_3 substrate as gate dielectric and a Cu plate as back-gate electrode. Carriers are depleted (accumulated) when the gate voltage V_G is negatively (positively) biased. (C) Temperature dependence of resistance in thin films with different Zn concentration x . The film thicknesses are 35 nm. (D) Electron density versus Zn concentration x of thin films compared with bulk single crystals. The sketch on the top depicts the topological phase transition of $(\text{Cd}_{1-x}\text{Zn}_x)_3\text{As}_2$ from a DSM to a trivial insulator. The critical doping level was estimated to be $x_c \sim 0.17$ (28). (E) Relationship of mobility and electron density of thin films and bulk samples (26).

$(\text{Cd}_{1-x}\text{Zn}_x)_3\text{As}_2$ alloy films are confirmed to have a (112)-oriented single phase, and their lattice constant agrees with Vegard's law (fig. S1) (25). We deposited the films through a stencil metal mask into a Hall bar shape with a typical channel width of 60 μm , and we used the SrTiO_3 substrates directly as gate dielectric in a back-gate configuration, as depicted in Fig. 1B.

As Zn is doped in Cd_3As_2 , the temperature dependence of the longitudinal resistance R_{xx} exhibits a transition from metallic to semiconducting behavior (Fig. 1C). Figure 1D shows electron density of the $(\text{Cd}_{1-x}\text{Zn}_x)_3\text{As}_2$ films as a function of the Zn concentration x along with the electron density reported for bulk single crystals for comparison (26). We observe a clear suppression of the residual electron density

as the Zn concentration increases. Notably, Zn doping restricts the electron density of thin films far more effectively than that of bulk single crystals, reflecting differences in growth methods (see also section S1). The electron mobility shown in Fig. 1E keeps values which exceed 10^4 cm^2/Vs for thin films, which is similar to those of bulk samples (26). This high mobility of the films in the low-electron density region makes it possible to investigate quantum transport phenomena down to the quantum limit.

Effective carrier suppression with less doping amount x is preferable in terms of keeping the Dirac dispersion of Cd_3As_2 . Because Zn_3As_2 is a semiconductor with ordinary band ordering, doping Zn into Cd_3As_2 not only compensates electrons but also lifts up the band inversion,

finally inducing a topological phase transition from a DSM to a trivial semiconductor. We predict the critical composition x_c for this transition around 0.17 based on the early magneto-optical measurements and theoretical studies (28) reporting that the energy gap of $(\text{Cd}_{1-x}\text{Zn}_x)_3\text{As}_2$ varies linearly with x and that the band inversion energy for Cd_3As_2 is 0.19 eV. Therefore, it is reasonable to conclude that at least the $x = 0.06$ and 0.11 samples are still in the band-inverted regime, as depicted in the schematics of Fig. 1D, and our magneto-transport data of the $x = 0.11$ sample presented later reflect the existence of a linear energy dispersion.

Figure 2A presents the magnetic field dependence of longitudinal resistance R_{xx} and Hall resistance R_{yx} . With the progressive suppression of the electron density, the amplitudes of the quantum oscillations are enhanced in R_{xx} and QH states start to appear with characteristic quantized plateaus in R_{yx} . Even in 35-nm-thick films, the observation of the QHE, which is essentially a 2D phenomenon, can be explained by the large Fermi velocity in the confinement direction. The sub-band splitting energy becomes as large as 80 meV in the case of the 35-nm-thick film in the $x = 0.11$ -doped sample, as estimated later. In such a situation,

only a few sub-bands are occupied by electrons at low carrier densities, and thus, the 2D nature of each sub-band emerges.

The observed Hall plateaus in $(\text{Cd}_{1-x}\text{Zn}_x)_3\text{As}_2$ films exhibit quantization at $\sigma_{xy} = sn \frac{e^2}{h} = v \frac{e^2}{h}$, where n is integers, s is the spin and valley degeneracy, and v is the filling factor, rather than at the half-integer-shifted values $\sigma_{xy} = s(n + \frac{1}{2}) \frac{e^2}{h}$ typically observed in gapless Dirac systems (29, 30). Lack of such a shift in σ_{xy} indicates the absence of the zero-energy $N = 0$ Landau level, which is expected based on the existence of a confinement gap in the sub-band (20). The energy dispersion on a momentum plane off the Dirac points is described well by a hyperbolic function with a finite gap at the bottom (Fig. 1A) (5), which has been resolved in the angle-resolved photoemission spectroscopy experiment of bulk Cd_3As_2 (6). In this sense, despite the confinement gap, the current system maintains the linear character of the energy dispersion and the observed QH states are explained well by Landau quantization of such a gapped Dirac dispersion.

Combined with field effect, it is possible to deplete the $(\text{Cd}_{1-x}\text{Zn}_x)_3\text{As}_2$ films further and even gate-tune them into a p-type regime, as indicated by a clear sign change in R_{yx} of the negatively biased $x = 0.17$ sample.

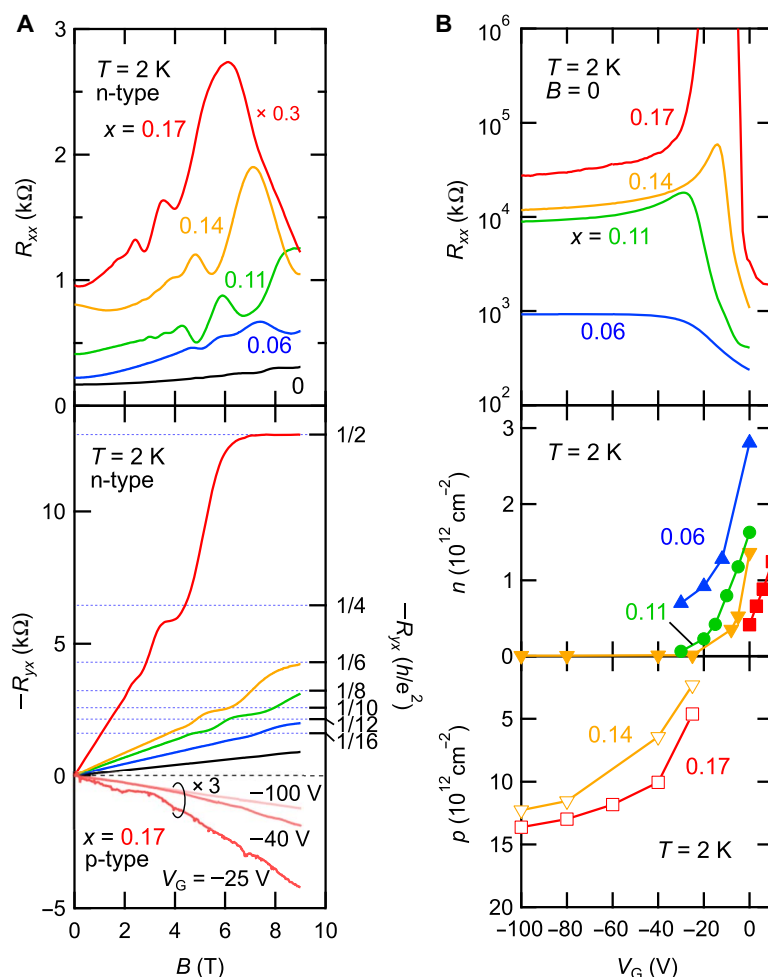


Fig. 2. Carrier density control in $(\text{Cd}_{1-x}\text{Zn}_x)_3\text{As}_2$ thin films. The film thicknesses are 35 nm. (A) Magnetic field dependence of longitudinal resistance R_{xx} and Hall resistance R_{yx} at 2 K. Quantized plateau regions appear in R_{yx} as the electron density is decreased. Applying a large negative gate voltage V_G induces a clear sign change of the Hall coefficient as shown for the case of $x = 0.17$. (B) V_G dependence of R_{xx} , sheet electron density n , and sheet hole density p at 2 K. We produce ambipolar behaviors by back-gating the samples in the case of $x \geq 0.11$. Accompanied by the carrier-type inversion, R_{xx} exhibits a peak that corresponds to the CNP. In the cases of $x = 0.11$ and 0.14, we observe multicarrier conduction with majority holes and minority electrons after reaching the CNP (see fig. S2).

Figure 2B presents the gate voltage V_G dependence of R_{xx} and the sheet carrier density. Following the carrier type inversion, a peak structure appears in R_{xx} , which corresponds to the charge neutrality point (CNP). The resistance value at the CNP exhibits a diverging behavior with x . The suppressed band inversion of Cd_3As_2 by Zn doping, and hence the increase of the energy gap in the confined sub-band, account for such a dependence. In the p-type regime for $x = 0.11$ and 0.14 , R_{yx} exhibits multicarrier conduction in contrast to single-hole conduction for $x = 0.17$ (fig. S2). The existence of residual electrons implies either the activation of electrons or the formation of electron-hole puddles (31). The deviation from the electron density's linear V_G dependence on the CNP may also be related to this inhomogeneity effect. On the other hand, we expect Zn doping to enhance the energy gap and hinder the puddle formation, eventually leading to single-hole conduction for $x = 0.17$. Both the quantum oscillations and the QH states are suppressed in the p-type region, probably due to the lower hole mobility or the heavier mass of the VB (fig. S2), which is consistent with previous reports (27, 32).

Gate modulation of QH states

Having established the carrier density control by field effect, we now investigate quantum transport at different Fermi-level positions. In Fig. 3, we summarize the results for $x = 0.11$, where the mobility is the highest among the samples. As we negatively increase V_G in -5 -V steps, Hall plateaus with smaller filling factors ν show up one

by one, and finally, the system reaches the quantum limit at $V_G = -20$ V, where all the electrons fall into the lowest Landau level for magnetic fields larger than $B = 8$ T. The magnetic field dependence of the QH states reveals that ν changes by 2 or 4 depending on the field, except for the last transition from $\nu = 2$ to 1. With the Fermi level located above the saddle point, the system only has a spin degeneracy of $s = 2$. Therefore, discrete jumps by 4 in ν , for example, around $B = 7$ T at $V_G = -5$ V, strongly suggest the existence of other sub-bands contributing to the QH states. To obtain a detailed picture of the Landau-level occupation and the resulting QH states when multiple sub-bands are involved, we must extract information about the band dispersion in $(\text{Cd}_{1-x}\text{Zn}_x)_3\text{As}_2$ films.

For this purpose, the oscillation terms of R_{xx} in lower fields are analyzed to estimate band properties such as Fermi momentum k_F , effective mass m^* , and Fermi velocity v_F . The Fermi momentum k_F and the Fermi surface area $S_F = \pi k_F^2$ are extracted from the periodicity of quantum oscillations using Fourier transformation. We note that the oscillations are mainly dominated by the single periodicity originating from the main sub-band (referred to as SB1), and the second occupied sub-band (referred to as SB2) is hardly resolved in the Fourier spectrum (figs. S3 and S4). On the other hand, we can estimate the effective mass m^* from the damping behavior of the oscillation amplitudes with temperature. Because the dielectric constant of the SrTiO_3 substrates heavily depends on temperature (33), we kept the electron density constant by tuning V_G at each temperature to investigate the

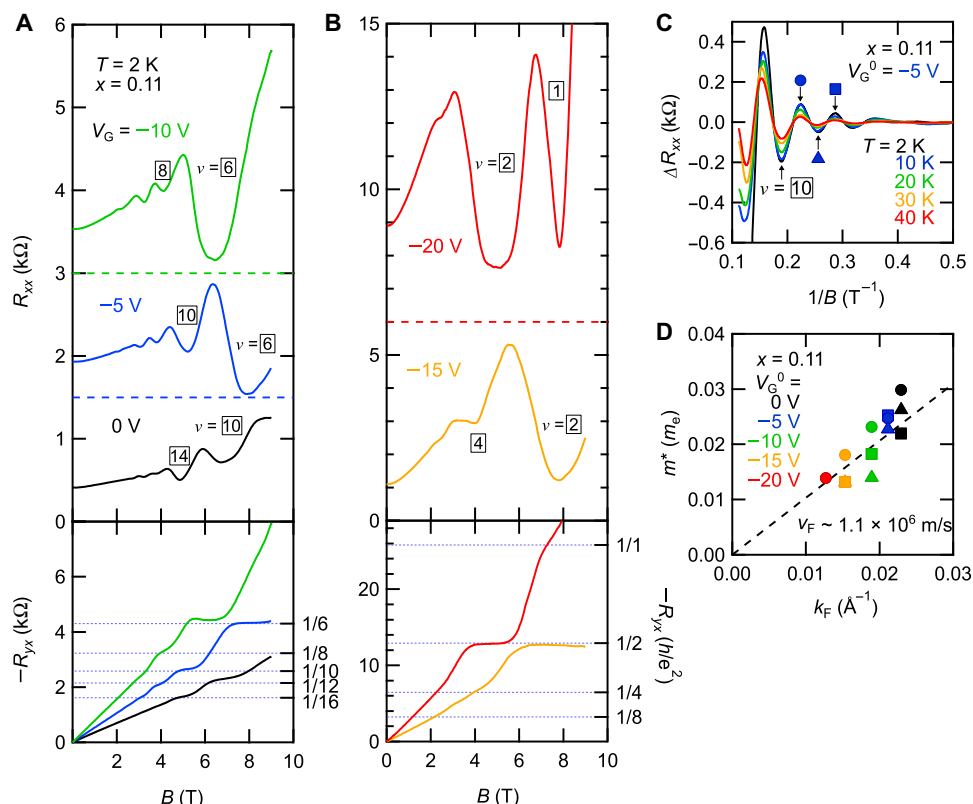


Fig. 3. Gate-tunable QH states in $(\text{Cd}_{1-x}\text{Zn}_x)_3\text{As}_2$ film ($x = 0.11$). (A and B) Magnetic field dependence of R_{xx} and R_{yx} measured at V_G from 0 to -20 V. The film thickness is 35 nm. QH states with smaller filling factors appear one by one as carriers are depleted. (C) Typical temperature-dependent oscillation terms of R_{xx} at lower fields. V_G^0 denotes V_G applied at 2 K, and V_G for higher temperatures is adjusted to keep the electron density constant. Here, the result for $V_G^0 = -5$ V is shown. (D) Relationship between the effective mass m^* and the Fermi momentum k_F of the main sub-band SB1. The different symbols represent m^* values estimated using different oscillation peaks and valleys as shown in (C) for $V_G^0 = -5$ V (see fig. S4A for the cases of other V_G^0 values). The Fermi velocity of SB1 is estimated to be about 1.1×10^6 m/s.

temperature dependence of the oscillations (fig. S3). Figure 3C shows a typical example of the oscillatory components of R_{xx} at different temperatures. Here, V_G^0 denotes the bias voltage at 2 K, and at higher temperatures, we adjusted V_G to keep the electron density of the sample constant, as deduced by low-field Hall measurements.

Figure 3D presents the change of m^* as a function of k_F for the main sub-band SB1. The effective mass exhibits a decreasing trend as the Fermi surface area shrinks. This behavior contrasts with the case of a parabolic band where m^* is independent of the Fermi surface area, suggesting that the energy dispersion remains linear for the $x = 0.11$ sample, as expected. We can ascribe the variation of the extracted effective mass at each Fermi level to the influence of the second occupied sub-band SB2, which may affect the appearance of oscillation amplitudes and bias the estimated m^* in either direction, underestimation or overestimation (see also section S4). By taking the average of these m^* , we estimate the Fermi velocity v_F of SB1 to be about 1.1×10^6 m/s, which is close to typical values previously reported for bulk Cd_3As_2 samples (13–15).

Analysis of Landau-level occupations

Having determined k_F of SB1, we can quantify the carrier occupation of SB2 as the difference between total electron density and partial electron density $n_{\text{SB1}} = k_F^2/2\pi$ of SB1. Figure 4A displays the electron density obtained from different experimental data as a function of V_G : n_{Hall} from the Hall coefficient shown in Fig. 2B, n_{QHE} from the quantization periodicity of the QH states shown in Fig. 4B, and n_{SB1} from k_F of SB1. Both n_{Hall} and n_{QHE} correspond to the total electron density of the sample and thus coincide over the entire range of V_G . On the other hand, n_{SB1} shows smaller values than the other two at $V_G > -15$ V, indicating the occupation of SB2. The occupation of the third sub-band (referred to as SB3) is also possible, but electrons in SB3 should be rapidly depleted under high magnetic fields. Therefore, for $V_G > -15$ V, it is reasonable to assume that two sets of Landau levels originating from SB1 and SB2 are involved in the emergence of the QH states.

Landau-level occupation within these multiple sub-bands can be well understood by plotting the filling factor ν against the inverse of

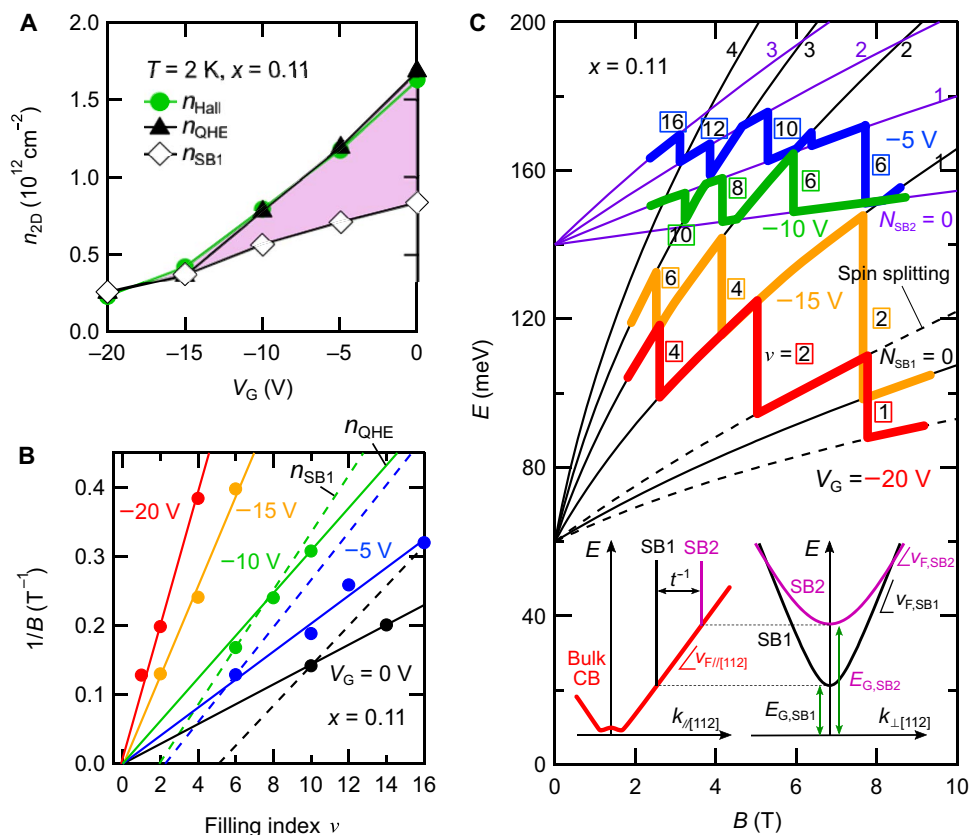


Fig. 4. Landau-level occupation identified from the QH states. (A) Gate voltage V_G dependence of carrier densities obtained from different experimental data. n_{Hall} is obtained from Hall measurements, n_{QHE} is obtained from the periodicity of the R_{xx} quantization, and n_{SB1} is obtained from the Fermi surface of the main sub-band SB1. While both n_{Hall} and n_{QHE} indicate the total electron density, n_{SB1} corresponds to the partial electron density in SB1. The difference between the total electron density and n_{SB1} at $V_G > -15$ V reflects the carrier population in the second sub-band SB2. (B) Filling factor plotted against the inverse of the magnetic field. The solid lines are fits to the data points with a fixed intercept at $\nu = 0$, yielding an estimate of n_{QHE} . The broken lines are drawn to pass each data point with Fermi surface area of SB1 as their slopes, representing the relationship between the filling factor and the magnetic field when only the presence of SB1 is assumed. We can use the intercept of the broken lines to estimate the Landau-level occupation in SB2 as explained in the main text (see also fig. S6). (C) Magnetic field dependence of Landau levels in SB1 and SB2, and their occupation transition at different Fermi levels. The thick lines overlaid on the Landau levels are the positions of chemical potential for each gate voltage. Accidental overlaps of the Landau levels in SB1 and SB2 result in jumps of the filling factor ν by 4. Inset is a sketch displaying the dispersions of the bulk CB along the confinement direction [112] (left), and those of sub-bands SB1 and SB2 perpendicular to [112] (right). We estimate each of the sub-band parameters to agree with the observed plateau transitions. The sub-band splitting energy ($E_{G,\text{SB2}} - E_{G,\text{SB1}}$) reaches about 80 meV at a confinement thickness $t = 35$ nm, corresponding to the Fermi velocity $v_{F,[112]}$ of about 7×10^5 m/s.

the magnetic field, as presented in Fig. 4B. Because each Landau level has a density of states of $\frac{eB}{h}$ at a given magnetic field B and total electron density n , quantized plateaus appear according to the relationship $n = \nu \frac{eB}{h}$. Thus, a linear fit to the data points yields the total electron density n_{QHE} as plotted by solid lines in Fig. 4A, regardless of the details of the multiple sub-bands (34). Next, we assume only the presence of SB1 and draw straight lines based on the Fermi surface area of SB1 from each data point, shown as broken lines in Fig. 4B (see also fig. S6). In contrast to the previous case, they do not pass through the origin but have a finite intercept on the horizontal axis. This is because we ignored the filled Landau levels and the corresponding density of states in SB2. From this viewpoint, it is reasonable to consider that the remaining intercept on the horizontal axis in return provides a rough estimation of the number of ignored Landau levels in SB2. For instance, the intercept of the broken line in the $V_G = -5$ V case is 2 for the $\nu = 6$ state, and this can be interpreted as the occupation of one Landau level from SB2 in addition to two Landau levels from SB1 to account for the $\nu = 6$ state. In this way, we can obtain the detailed breakdown of the filling factors between SB1 and SB2 at each magnetic field (fig. S6).

Figure 4C summarizes the magnetic field dependence of the Landau-level occupation at each gate voltage. By applying the semiclassical quantization condition to the hyperbolic dispersion (see also section S6), the Landau-level energies of confined Cd_3As_2 with a hyperbolic dispersion can be formulated as follows

$$E_N = \hbar v_F \sqrt{\frac{2e(N + \gamma)B}{\hbar} + \left(\frac{E_G}{\hbar v_F}\right)^2}$$

Here, N is the Landau index and v_F is the Fermi velocity in the in-plane direction ($\perp[112]$). γ reflects the Berry phase Φ and is defined as $\gamma = \frac{1}{2} - \frac{\Phi}{2\pi}$. We simply assume $\gamma = 0.5$ for the trivial case because of the gapped dispersion. E_G is the energy difference between the bulk Dirac points and the bottom of each sub-band, corresponding to the confinement-induced gap. On the basis of the information of electron density, the number of occupied Landau levels in SB1 and SB2, and field positions of the plateau transitions, we estimate the following parameters to agree with the experimental observations: $E_{G, \text{SB1}} = 60$ meV, $E_{G, \text{SB2}} = 140$ meV, $v_{F, \text{SB1}} = 1.1 \times 10^6$ m/s as determined in Fig. 3D, and $v_{F, \text{SB2}} = 8 \times 10^5$ m/s (see section S6 for a more detailed description of the estimation procedure). The occupation of SB2 above $E = 140$ meV opens the possibility that the Landau levels overlap at the Fermi level and hence certain filling factors are skipped. To account for the plateau transition from $\nu = 2$ to $\nu = 1$, spin-split levels are calculated by introducing the Zeeman term $\mu_B g B$, with $g = 25$ (12). Moreover, on the basis of Fig. 4C, the sub-band splitting energy ($E_{G, \text{SB2}} - E_{G, \text{SB1}}$) reaches 80 meV. Taking into account the confinement thickness of 35 nm, we can estimate the Fermi velocity in the confinement direction ($v_{F/[112]}$) to be around 7×10^5 m/s. Thus, combined with the systematic Fermi-level control, the band structure of the original 3D bulk is also accessible through the comprehensive analysis of QH states in 2D confined films.

DISCUSSION

Here, we assume that the sub-band splitting energy is kept constant over the investigated range of gate voltage and magnetic field for quantitative analysis, and it captures the essential picture of the experimentally ob-

served QH states well. For more strict quantification, a self-consistent calculation is preferable because the possible charge transfer between the spatially separated wave functions of the sub-bands can also induce a shift of sub-band splitting energy (35). Compared to the GaAs/ $\text{Al}_x\text{Ga}_{1-x}\text{As}$ wide-well case (35), however, we note that the possible shift (approximately a few millielectron volts) only has negligible effects because of the large confinement-induced sub-band splitting (80 meV) in the Dirac dispersion of $(\text{Cd}_{1-x}\text{Zn}_x)_3\text{As}_2$.

In addition to the detailed analysis of the QH states, systematic chemical doping and carrier control in thin films of Cd_3As_2 provide new opportunities to investigate the exotic phenomena predicted in the DSMs. As the Zn concentration dependence of the peak resistance at the CNP in Fig. 2B reveals the suppression of band inversion in Cd_3As_2 , Zn doping is useful in modulating the band structure of DSMs. In particular, the controllability of the Dirac point positions in the momentum space by Zn doping provides an additional controlling parameter for investigating the topological transport phenomena such as chiral anomaly, spin Hall effect (36), Fermi arc-mediated quantized transport (37), and potential topological superconductivity (38, 39). The electrostatic carrier control is also useful for investigating these phenomena because they appear dominantly when the Fermi level is close to the Dirac points.

Whereas Zn doping induces the topological phase transition from 3D DSMs to a trivial insulator, phase transitions into other 2D quantum phases such as quantum spin Hall insulator and quantum anomalous Hall insulator are possible in quantum wells and heterostructures (5, 20, 40). For the detection of these quantum phases, tuning the Fermi-level position into the bulk gap to single out the edge transport is unavoidable. In this sense, our findings on the quantitative analysis of the confinement effect and the carrier control across the CNP also lay the foundation for these attempts.

In summary, we have studied QH states of 2D confined Cd_3As_2 films where the carrier density is systematically controlled by chemical doping with Zn and electrostatic gating. Benefiting from comprehensive analysis down to the quantum limit, Landau-level formation from the linearly dispersed multiple sub-bands and their occupation transitions depending on the Fermi-level position have appeared. Our work on the chemical doping and Fermi-level tuning for Cd_3As_2 thin films also provides an important basis for the pursuit of transport manifestation of 3D Dirac Fermions and topological phase transitions predicted to occur in DSMs.

MATERIALS AND METHODS

Film growth

For the growth of $(\text{Cd}_{1-x}\text{Zn}_x)_3\text{As}_2$ thin films, Cd_3As_2 and Zn_3As_2 layers with a designed thickness ratio were first deposited in a stacking manner on SrTiO_3 substrates, which were then capped with TiO_2 (1 nm)/ MgO (5 nm)/ Si_3N_4 (200 nm). A patterned stencil metal mask was set on the substrate during the deposition process to make the film into the Hall bar shape. The TiO_2 and MgO capping layers were also deposited successively through the mask, but it was retracted when depositing the Si_3N_4 layer so that the latter covered the entire film, including the Hall bar edges. All the layers were deposited using a pulsed laser deposition technique at room temperature. The films were then subjected to annealing at 600°C in air. This annealing process promoted crystallization of the films and the formation of a solid solution of the two arsenide layers. In the presence of the capping materials, temperature could be elevated without evaporation of the deposited arsenide films or

any undesired chemical reactions taking place, resulting in $(\text{Cd}_{1-x}\text{Zn}_x)_3\text{As}_2$ alloy films of high crystalline quality. The Zn concentration was simply controlled by modifying the thickness ratio of the deposited Cd_3As_2 and Zn_3As_2 layers.

Transport measurements

Low-temperature transport measurements were carried out in a Physical Property Measurement System (Quantum Design) with a magnetic field of up to 9 T and at a base temperature of 2 K. Hall measurements were conducted using a lock-in technique. The excitation current was kept constant at 0.1 μA with a frequency of 13 Hz. For field effect measurements, dc bias voltage was applied to the Cu gate electrode.

SUPPLEMENTARY MATERIALS

Supplementary material for this article is available at <http://advances.sciencemag.org/cgi/content/full/4/5/eaar5668/DC1>

section S1. X-ray diffraction and estimation of Zn concentration x
 section S2. Magnetotransport of $(\text{Cd}_{1-x}\text{Zn}_x)_3\text{As}_2$ films in p-type region
 section S3. Temperature dependence of magnetotransport
 section S4. Estimation of Fermi momentum and effective mass
 section S5. Analysis of Landau-level occupation
 section S6. Estimation of sub-band parameters
 fig. S1. X-ray diffraction and estimation of Zn concentration x .
 fig. S2. Magnetotransport of $(\text{Cd}_{1-x}\text{Zn}_x)_3\text{As}_2$ films in p-type region.
 fig. S3. Temperature dependence of magnetotransport.
 fig. S4. Estimation of Fermi momentum k_F .
 fig. S5. Estimation of effective mass m^* .
 fig. S6. Analysis of Landau-level occupation.
 fig. S7. Appearance of QH states in the samples with different thicknesses.
 fig. S8. Determination of field positions for Hall plateau transitions.
 fig. S9. Estimation of sub-band parameters.
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