

Giant nonreciprocal emission of spin waves in Ta/Py bilayers

Jae Hyun Kwon,^{1*} Jungbum Yoon,^{1*} Praveen Deorani,¹ Jong Min Lee,¹ Jaivardhan Sinha,² Kyung-Jin Lee,^{3,4} Masamitsu Hayashi,^{2,5} Hyunsoo Yang^{1†}

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Spin waves are propagating disturbances in the magnetization of magnetic materials. One of their interesting properties is nonreciprocity, exhibiting that their amplitude depends on the magnetization direction. Nonreciprocity in spin waves is of great interest in both fundamental science and applications because it offers an extra knob to control the flow of waves for the technological fields of logics and switch applications. We show a high nonreciprocity in spin waves from Ta/Py bilayer systems with out-of-plane magnetic fields. The nonreciprocity depends on the thickness of Ta underlayer, which is found to induce an interfacial anisotropy. The origin of observed high nonreciprocity is twofold: different polarities of the in-plane magnetization due to different angles of canted out-of-plane anisotropy and the spin pumping effect at the Ta/Py interface. Our findings provide an opportunity to engineer highly efficient, nonreciprocal spin wave-based applications, such as nonreciprocal microwave devices, magnonic logic gates, and information transports.

INTRODUCTION

The unidirectional propagation of signals is the key to modern electronic logic circuits and fundamental applications, such as diodes, isolators, gyrators, and circulators (1). Because nonreciprocity offers an opportunity to control the wave devices flowing only to a desirable direction, the nonreciprocity in waves has been extensively studied, and wave devices with nonlinear propagating properties are introduced in the fields of microwaves (2), acoustics (3, 4), and photonics (5, 6). A typical nonreciprocity ratio of nonreciprocal microwave devices (7) is $\sim 10^3$, and the nonreciprocity of acoustics (3) and photonics (5) is reported to be $\sim 10^4$.

In rapidly developing areas of magnonics, the spin waves (8–14) can be the information carrier for magnonic logic circuits (15–19). The specific property of the nonreciprocity might be a benefit for the enhancement of logic circuits based on magnetic waveguide. Therefore, the characterization and tuning of the nonreciprocity in magnonics have been intensively investigated (20–31), including the relation between the nonreciprocity and spin pumping (32) or Dzyaloshinskii-Moriya interaction (DMI) (33, 34). The nonreciprocity, due to its intrinsic characteristic in the magnetostatic surface wave (MSSW) mode (35), together with Joule-heat-free transport in gigahertz frequencies (30) and compatibility with the spin Hall current system (36, 37) makes spin wave one of the promising candidates for future data carriers.

However, the suitability of nonreciprocal spin waves for the above applications is limited to a small nonreciprocity parameter, κ . For example, it was reported that a typical value of κ in MSSW is 0.6 to 0.9 in a permalloy (Py) strip by using time-resolved spin wave spectroscopy (see fig. S1) (23, 26) and 4 to 5 by using spatially resolved Brillouin light-scattering spectroscopy (22). The value of κ in an yttrium iron garnet film, which has a low saturation magnetization, is higher compared to that in Py (27). However, a large value of κ in a typical 3d ferromagnet, such as

Py, is essential for the realization of spin wave-based applications, because Py is compatible with the prevailing microfabrication technology.

Here, we report giant nonreciprocal emission of spin waves in a Ta/Py bilayer with out-of-plane bias fields, with κ values up to 14 and 60 in time and frequency domains, respectively. An optical image of the spin wave device used in this study is shown in Fig. 1A. A pulse inductive technique was used for measurements, in which an electric pulse was applied to the right asymmetric coplanar strip (ACPS) for the excitation of the spin wave in a ferromagnetic thin film (20-nm Py) on top of various thicknesses (t_{Ta}) of Ta from 1.2 to 8.9 nm. The spin wave propagating along the x direction is detected using a sampling oscilloscope at the left ACPS. An out-of-plane bias field ($|H_z| < 490$ mT) is applied to the sample during the measurements. All measurements were performed at room temperature.

The nonreciprocity in MSSW is due to either the interference between the spin waves produced by the x and z components of rf local excitation fields (antenna-spin wave coupling) or the local concentration of the spin waves at the top and bottom surfaces of the magnetic film. In our case, the amplitude of spin waves across the Py thickness can be assumed to be uniform because $kt_{\text{Py}} \ll 1$ (22, 35), where k ($2\pi/\lambda$, ~ 0.63 rad/ μm ; λ is the wavelength of spin waves) is the wave vector and t_{Py} (20 nm) is the thickness of Py. Therefore, the basic mechanism of nonreciprocity in our system is the antenna-spin wave coupling. In addition, we discuss two reasons for the observed high nonreciprocity arising at the Ta/Py interface. The first reason is the tilted out-of-plane anisotropy induced by the Ta underlayer, and the second one is the spin pumping effect at the Ta/Py interface. The vibrating sample magnetometry (VSM) and anisotropic interface magnetoresistance (AIMR) measurements indeed show the presence of an interfacial anisotropy at the Ta/Py interface. Taking this anisotropy into account, the micromagnetic simulations support the observed experimental results.

RESULTS

Figure 1B shows the contour plot of the field-dependent measurements in the time domain for the device with $t_{\text{Ta}} = 6.1$ nm. A representative

¹Department of Electrical and Computer Engineering and National University of Singapore Nanoscience and Nanotechnology Initiative, National University of Singapore, Singapore 117576, Singapore. ²National Institute for Materials Science, Tsukuba 305-0047, Japan. ³Department of Materials Science and Engineering, Korea University, Seoul 02841, Korea. ⁴KU-KIST Graduate School of Converging Science and Technology, Korea University, Seoul 02841, Korea. ⁵Department of Physics, The University of Tokyo, Bunkyo, Tokyo 113-0033, Japan.

*These authors contributed equally to this work.

†Corresponding author. Email: eleyang@nus.edu.sg

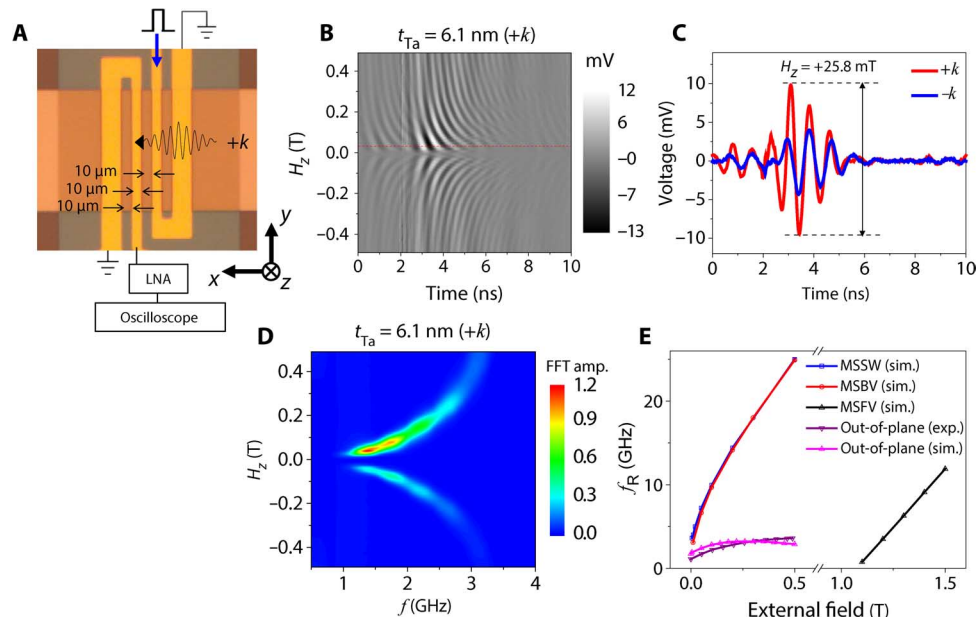


Fig. 1. Device structure and spin wave signals with out-of-plane bias fields. (A) The device consists of Si/SiO₂ substrate/Ta (t_{Ta} nm)/Py (20 nm)/Ru (3 nm)/SiO₂ (35 nm). The thickness of the Ta underlayer ranges from 1.2 to 8.9 nm. ACPs are used to apply local excitation fields and inductively detect magnetization precessions using a sampling oscilloscope after amplification with a low-noise amplifier (LNA). (B) Contour plot of spin waves in the time domain with various H_z for $t_{\text{Ta}} = 6.1$ nm. The fields ($|H_z| < 490$ mT) are applied perpendicularly to the film plane and are smaller than the demagnetization field, and thus cannot excite MSFV mode. (C) Spin wave packet at $H_z = 25.8$ mT (corresponding to the red dotted line in Fig. 1B). The amplitude (amp.) is defined as the difference between the maximum and minimum of the wave packet. (D) Contour plot of FFT of time domain signals. The spin wave resonance frequency (f_R) increases with increasing $|H_z|$. (E) The range of f_R corresponding to the external fields differs from that of the conventional spin wave modes, such as MSSW, MSBV, and MSFV. The f_R plots for MSSW, MSBV, MSFV, and out-of-plane (sim.) are obtained from the simulations, in which the material parameters of Py are used. The experimental data of f_R for the out-of-plane (exp.) are obtained from the sample of $t_{\text{Ta}} = 6.1$ nm.

spin wave packet measured at $H_z = 25.8$ mT is shown in Fig. 1C for different propagation directions ($\pm k$). The shape of this wave packet closely resembles that of the MSSW. The fast Fourier transform (FFT) of time domain signals in Fig. 1D shows that the traveling spin wave resonance frequency (f_R) increases with increasing $|H_z|$ for the device with $t_{\text{Ta}} = 6.1$ nm. Figure 1D also shows that the FFT intensity of spin wave packets is stronger for positive H_z , as compared to that for negative H_z , demonstrating nonreciprocal emission of spin waves depending on the polarity of the applied out-of-plane fields. For a comparison, the MSSW (in-plane bias fields), the magnetostatic backward volume (MSBV), the magnetostatic forward volume (MSFV), and the MSSW-like (out-of-plane fields) mode of this study are shown in Fig. 1E.

The amplitude of wave packets for devices with different t_{Ta} is shown as a color plot in Fig. 2A. The maximum value of the spin wave amplitude was observed at $t_{\text{Ta}} = 6.1$ nm for the positive H_z and at $t_{\text{Ta}} = 3.3$ nm for the negative H_z . Figure 2B shows that the amplitude at $H_z = -25.8$ mT decreases as t_{Ta} increases, whereas the amplitude at $H_z = +25.8$ mT initially increases (up to $t_{\text{Ta}} = 6.8$ nm) and then decreases, as summarized in Fig. 2C. The first regime ($t_{\text{Ta}} < 6.8$ nm) is explained by the tilted anisotropy induced by the Ta layer, and the second regime ($t_{\text{Ta}} \geq 6.8$ nm) is attributed to the spin pumping effect, as discussed later. The spin wave nonreciprocity (κ) is defined as the ratio of spin wave amplitudes at positive and negative fields; $\kappa = A(+H_z)/A(-H_z)$. The maximum value of κ in the time domain is ~ 14 for the device with $t_{\text{Ta}} = 8.2$ nm, as shown in the inset of Fig. 2C, and the κ reaches up to 60 in the frequency domain

(the details of κ in the time and frequency domains are documented in sections S2 and S3).

We propose that the effect of the Ta underlayer is to induce a canted interfacial magnetic anisotropy, which affects the nonreciprocity factor. To confirm the presence of interfacial anisotropy induced by Ta underlayers, we conducted the out-of-plane AIMR measurements on all the devices, and the results are shown in fig. S4C. Because of AIMR, the in-plane electrical resistance in the perpendicular anisotropy materials depends on the out-of-plane component of the magnetization (38). In Fig. 3, ΔR_{OP} at $+H_z$ and $-H_z$ are plotted versus the Ta underlayer thickness, where ΔR_{OP} is the maximum change in the resistance of the device when the magnetic field is swept along the z direction. ΔR_{OP} is a measure of AIMR that originates from the Ta/Py interface, due to anisotropic magnetic scattering of electrons (38). ΔR_{OP} changes its polarity across $t_{\text{Ta}} \sim 3$ nm, which suggests a change in anisotropy direction in the devices. It is consistent with the trend observed in a polarity change of spin wave amplitudes in Fig. 2B.

For further confirmation of the presence of induced anisotropy in our sample, VSM measurements were conducted, and a weak perpendicular anisotropy in the z direction was indeed observed. The anisotropy was verified from an inner loop from VSM hysteresis as shown in fig. S5, which is in line with a recent report (39). The z component of the magnetization shows a finite remanence value (M_r) at $H_z = 0$. The measured M_r with rotating the sample around the x axis shows the minimum at a positive angle for $t_{\text{Ta}} = 1$ nm and at negative angles for $t_{\text{Ta}} = 3, 5,$ and 10 nm,

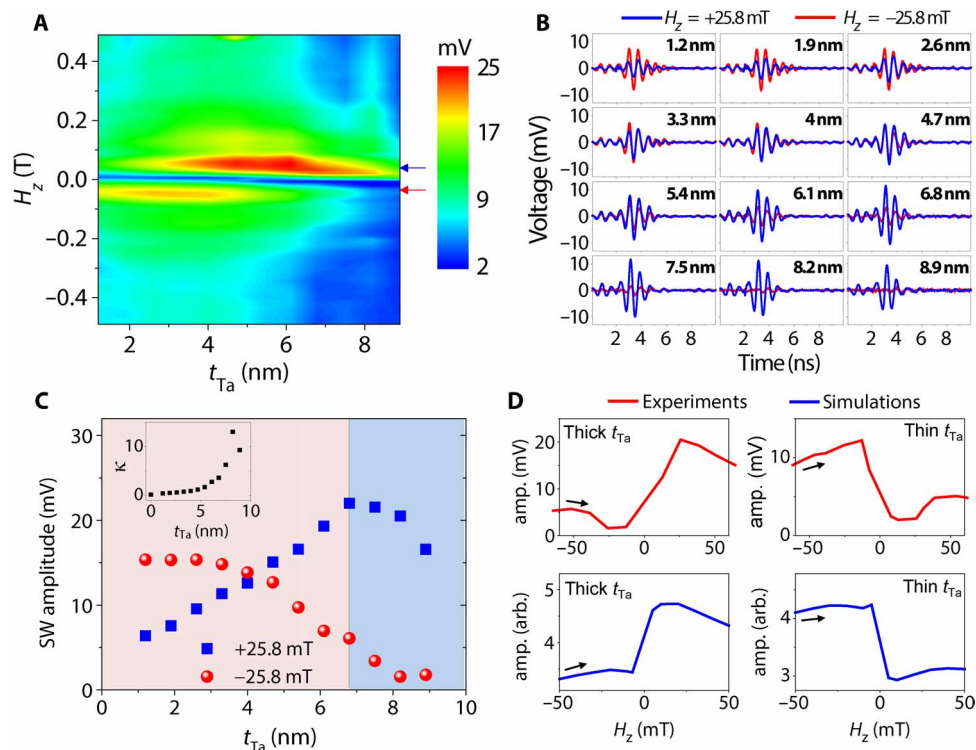


Fig. 2. Change in spin wave amplitudes with different t_{Ta} . (A) Contour plot of spin wave (SW) amplitudes as a function of t_{Ta} and H_z . The red and blue arrows on the right side indicate $H_z = \pm 25.8$ mT. (B) Spin wave packets at $H_z = \pm 25.8$ mT. The amplitude at $-H_z$ is higher than that at $+H_z$ in the device for $1.2 < t_{Ta} < 3.3$ nm. The amplitude at $+H_z$ is higher than that at $-H_z$ for $4.7 < t_{Ta} < 8.9$ nm. (C) Amplitudes of the wave packets from Fig. 2B are plotted with different t_{Ta} , and the nonreciprocity (κ) is shown in the inset. κ is found to be ~ 0.2 for $t_{Ta} = 0$ nm. (D) Simulation results are shown for comparison with experimental results. H_z sweeps from negative to positive fields. The simulations show good agreement with experimental results. The experimental data are obtained from a thin case ($t_{Ta} = 1.9$ nm) and a thick case ($t_{Ta} = 8.2$ nm). arb., arbitrary units.

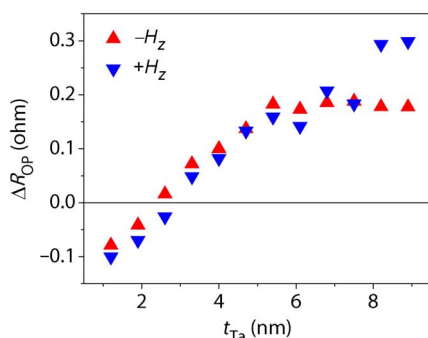


Fig. 3. Anisotropic interface magnetoresistance. ΔR_{OP} is plotted with different t_{Ta} and shows a polarity change at $t_{Ta} \sim 3$ nm and an increase with increasing t_{Ta} . It indicates that different t_{Ta} changes the interfacial anisotropy directions.

indicating a change in the canted angle of perpendicular anisotropy for different t_{Ta} . For the thicker t_{Ta} , the uniaxial anisotropy is canted to the y axis with the projected M_y along the $-y$ direction with $H_z > 0$, whereas the M_y is along the $+y$ direction for the thinner t_{Ta} with $H_z > 0$.

To better understand the relation between anisotropy and observed nonreciprocity, we performed micromagnetic simulations (see fig. S6). Simulations were carried out on a bilayer structure, with the bottom

layer having a perpendicular uniaxial anisotropy. The spin wave amplitude obtained from simulations shows a qualitative agreement with the measured data as shown in Fig. 2D. From the investigation of magnetization configurations in simulated results, we found that the spin wave nonreciprocity is correlated with the polarity of the y component of the magnetization. Because the magnetization has the y component, the observed spin waves are MSSW-like. The tilting direction of the aforementioned anisotropy depends on the t_{Ta} at fixed $\pm H_z$. When the applied H_z varies, the y component of the magnetization changes its sign, leading to a change in relative amplitude of spin waves at $\pm H_z$ (21–23, 31). In addition, for a better understanding of the antenna–spin wave coupling, we spatially and temporally analyzed our micromagnetic simulation results with the relation between the excitation magnetic field from the antenna and the local magnetization at the bottom of the antenna (see fig. S7).

Although a change in the anisotropy direction depending on t_{Ta} can explain the trend of the nonreciprocity, a high nonreciprocal factor especially at thicker Ta ($t_{Ta} \geq 6.8$ nm) cannot be fully explained using the above model. This discrepancy between the results from the micromagnetic simulations and the experimental observation suggests that a mechanism other than the interfacial anisotropy is also in operation and becomes more dominant at thicker Ta layers. To further understand the origin of high nonreciprocity, the loss mechanisms for spin wave energy should be considered. Apart from the usual energy loss mechanisms in a ferromagnet, which are quantified by the Gilbert

damping α , a heavy metal Ta underlayer can also absorb the energy from magnetization dynamics via spin pumping (32). Spin pumping is the process in which the magnetization precession in a ferromagnet (Py) generates spin current density $\vec{J}_S \propto \vec{m} \times (\partial \vec{m} / \partial t)$ into the adjacent paramagnetic layer (Ta) (40–42). Because of high spin-orbit coupling in Ta, this spin current is converted to a transverse charge current $\vec{J}_{C \propto} \vec{J}_S \times \vec{m}$ via inverse spin Hall effect (ISHE), which can be measured as a transverse electromotive force across the device.

We performed spin pumping–induced ISHE measurements to investigate energy loss from spin waves in Py to Ta. A schematic of the spin pumping measurement is shown in Fig. 4A. In the geometry of our experiment, a spin pumping–induced ISHE signal (V_{sp}) can be observed only if there exists a nonvanishing y component of the magnetization at the interface, which accounts for the observation of MSSW-like wave packets in Fig. 2B. The ISHE signals measured at an *rf* frequency of 1.3 GHz corresponding to the spin wave resonance frequency at $|H_z| = 25.8$ mT are shown in Fig. 4B. The amplitude of spin pumping signals exhibits an abrupt jump at $t_{Ta} = 6.8$ nm as shown in Fig. 4C, which can be related to the spin diffusion length in Ta (~ 7 nm) (43). The amplitude of V_{sp} is a measure of loss of the total spin wave energy that is absorbed into the Ta layer (44). Thus, above $t_{Ta} \sim 6.8$ nm, the spin wave amplitude should decrease because of the strong spin pumping effect. In Fig. 2C, this decrease in the spin wave amplitude is clearly observable for both positive and negative applied fields for $t_{Ta} \geq 6.8$ nm.

If the spin pumping changes damping equally for the $+H_z$ and $-H_z$ spin waves, both wave amplitudes will be reduced proportionately, leading the nonreciprocity to be unchanged. Therefore, the spin pumping must unequally affect the spin waves for two field directions ($\pm H_z$), to explain the observed giant nonreciprocity, especially at thicker

Ta ($t_{Ta} \geq 6.8$ nm). We propose a plausible mechanism of the spin wave amplitude–dependent damping, in which the spin pumping–induced Gilbert damping of large-amplitude spin waves at $+H_z$ is smaller than that of small-amplitude spin waves at $-H_z$ (see section S8). Consequently, a large-amplitude (small-amplitude) spin wave experiences less (more) attenuation when propagating, resulting in an increased nonreciprocity at thicker Ta.

To confirm the spin pumping effect, we performed the ferromagnetic resonance (FMR) measurements with the same devices of the spin pumping measurements. Taking the spin pumping effect into account, the presence of Ta layer enhances the damping parameter of Py above the intrinsic value (45, 46)

$$\alpha_{\text{eff}} = \alpha_0 + \Delta\alpha = \alpha_0 + \frac{g\mu_B}{4\pi M_s t_{\text{FM}}} g_{\text{eff}}^{\uparrow\downarrow}$$

where α_0 is the intrinsic damping, $\Delta\alpha$ is an additional damping because of the spin pumping effect, g (2.2) is the g -factor, μ_B is the Bohr magneton, t_{FM} (20 nm) is the thickness of the Py film, M_s is the saturation magnetization, and $g_{\text{eff}}^{\uparrow\downarrow}$ is the effective spin-mixing conductance. We extracted the additional damping constant ($\Delta\alpha$) from the FMR spectra as shown in Fig. 4C. The plotted damping enhancement as a function of the Ta thickness shows good agreement with spin pumping signals.

DISCUSSION

From the first principles calculation, the $5d$ orbital of Ta, which has a large spin-orbit coupling, is strongly hybridized with the $3d$ orbital of Fe,

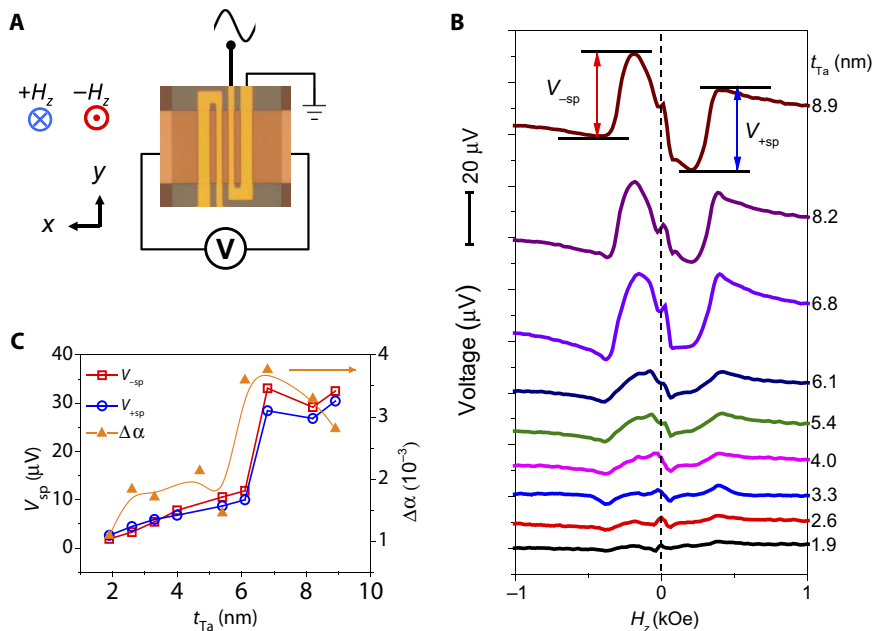


Fig. 4. Spin pumping measurements by spin waves and ferromagnetic resonance data with out-of-plane H_z . (A) The spin pumping voltage is measured at the end of both electrodes. The applied 1.3-GHz microwave continuously excites local magnetization oscillations. (B) The amplitude of spin pumping signals increases with thicker t_{Ta} . It is expected that the amplitude shows the maximum for $t_{Ta} \sim$ spin diffusion length (7 nm) in Ta. The spin pumping signal with $t_{Ta} = 1.2$ nm is used as a background signal and subtracted from that of all other devices. (C) A sudden increase in V_{sp} is observed at $t_{Ta} = 6.8$ nm. It is inferred that the spin pumping effect plays a significant role for $t_{Ta} \geq 6.8$ nm. The plotted damping enhancement ($\Delta\alpha$) as a function of the Ta thickness shows good agreement with spin pumping signals.

giving rise to strong perpendicular magnetic anisotropy at the interface (47). The anisotropy properties, such as an axis or strength, are sensitive to intermixing, roughness, oxidation, and hybridization at the interface, which are affected by different t_{Ta} . The origin of the canted perpendicular anisotropy to the y axis could be a change in those parameters depending on t_{Ta} . The canted anisotropy only occurs at the Ta/Py interface, and the anisotropy direction becomes closer to the x axis away from the interface because of the shape anisotropy. It has been reported earlier that a Ru seed layer on the top of NiFe can induce surface/interface anisotropy (48), and Ta or Ru underlayer with CoFeB can result in a perpendicular anisotropy (49). In our experiment, the nonreciprocity is found to change with different t_{Ta} , whereas the thickness of Ru is kept constant at 3 nm. Therefore, we conclude that the Ta underlayer induces a progressive change in the spin wave properties.

Finally, we investigated the possible role of DMI in our devices. DMI is observed in a broken inversion symmetry (50, 51) and at the interface or surface of magnetic multilayers (52, 53) and could be expected to have an opposite effect in the inverted sample structure. Therefore, we performed spin wave measurements in devices with the Ta layer on the top of the Py layer (see fig. S9). It turns out that the nonreciprocal ratio in the inverted devices show a similar trend as that of Ta underlayer structures. In our case, we conclude that the structural sequence between the Ta and Py layers does not affect the nonreciprocal emission of spin waves.

It must be noted that our observations are different from the earlier reports on spin wave nonreciprocity studied only in the MSSW mode (22, 23, 26, 28). To compare our results with the MSSW mode, we have also measured the nonreciprocity value with sweeping in-plane magnetic field, H_y . It is found that a small value (~ 0.7) of nonreciprocity is observed regardless of the Ta thickness, similar to previous reports (54) (see fig. S1).

The observed very large nonreciprocity of spin waves in the Ta/Py bilayer can be attributed to the combined effects, such as the polarity change in the in-plane magnetization due to the canted out-of-plane anisotropy and the spin pumping process at the Ta/Py interface. Utilizing an intimate coupling between the interfacial spin configuration and spin pumping effect, even larger nonreciprocity can be achieved. In addition, noncollinear spin configuration due to DMI may enhance the nonreciprocity ratio (33). Our intriguing findings will be useful for future spin wave-based applications, such as spin wave information transports and logic devices because of the high nonreciprocal ratio.

MATERIALS AND METHODS

Sample preparation

Si/SiO₂ substrates/Ta (t_{Ta} nm)/Ni₈₁Fe₁₉ (20 nm)/Ru (3 nm) was patterned in a rectangular shape (120 × 340 μm²). Magnetic fields were not applied during the sputtering. A wafer was being rotated during sputtering to avoid a field deposition under any stray field effects from the magnetron sputtering gun or the chamber. The SiO₂ (35 nm) insulating layer was deposited to isolate an ACPS from the Py layer. The ACPS consists of Ta (5 nm)/Au (100 nm). The signal line width is 10 μm. The gap (edge to edge) between the signal line and the ground conductor is 10 μm, and the distance between the signal lines of two ACPSs is 10 μm as shown in Fig. 1A.

Measurement setup

We used pulsed inductive microwave magnetometry to generate and measure the spin wave signals. Spin waves were generated by applying

a 1.8-V electric pulse to the excitation ACPS to produce local rf magnetic fields around the signal line. The rise and fall times of the electric pulse were ~ 70 and ~ 80 ps, respectively. The pulse width was ~ 100 ps, and the pulse cycle frequency was 100 MHz. The excited spin waves propagated through the Py strip and were inductively detected by the detection ACPS (8, 55). The detected spin wave signals were then amplified using a low-noise amplifier (29-dB gain) and subsequently passed into a sampling oscilloscope. The waveforms were averaged 10,000 times by the oscilloscope to improve the signal-to-noise ratio, followed by subtraction of reference signals to obtain the pure wave packets. To obtain the additional damping constant, vector network analyzer-based FMR measurements were performed by a frequency-swept technique with the same devices used in the spin pumping measurements.

Simulations

We used Object-Oriented Micromagnetic Framework developed by M. J. Donahue and D. G. Porter (see <http://math.nist.gov/oommf/>) to simulate spin waves by solving the Landau-Lifshitz-Gilbert equation. The magnetic film structure used in our simulations consists of two Py layers (top and bottom layer) with different anisotropy axes. The anisotropy axis was tilted 0.4° away from the z axis toward the $-y$ direction for thick t_{Ta} and toward the $+y$ direction for thin t_{Ta} . The value of the uniaxial anisotropy constant (K_u), saturation magnetization (M_s), and exchange stiffness constant (A) were 7.0×10^5 J/m³, 8.0×10^5 A/m, and 1.3×10^{-11} J/m, respectively. The Gilbert damping constant (α) and the gyromagnetic ratio (γ) in our simulations were set to 0.01 and 2.32×10^5 m/(A·s), respectively. The dimensions of the top and bottom layers were $8000 \times 200 \times 16$ nm³ and $8000 \times 200 \times 4$ nm³, respectively, and the simulation cell size was $4 \times 200 \times 4$ nm³. To excite spin wave dynamics, a magnetic field pulse of amplitude 3 mT is applied to a volume of $200 \times 200 \times 20$ nm³ at a distance of 1 μm from one end of the Py strip, and the propagating spin waves were detected by the summing of the dynamic magnetization over a volume of $800 \times 200 \times 20$ nm³ at 5 μm away from the excitation source. The out-of-plane bias field was varied from -500 to $+500$ mT during our simulations.

SUPPLEMENTARY MATERIALS

Supplementary material for this article is available at <http://advances.sciencemag.org/cgi/content/full/2/7/e1501892/DC1>

- section S1. Nonreciprocal ratios of surface spin wave mode with in-plane fields
- section S2. Nonreciprocity ratios from spin wave amplitudes in the time domain
- section S3. Nonreciprocity from spin wave intensity in the frequency domain
- section S4. Anisotropic magnetoresistance measurements
- section S5. VSM measurements
- section S6. Simulations: Out-of-plane hysteresis and nonreciprocity
- section S7. Antenna-spin wave coupling
- section S8. Spin wave amplitude-dependent damping
- section S9. Ta on the top of the Py layer
- section S10. Dependence of nonreciprocity on different directions of field sweep
- section S11. Dependence of nonreciprocity on the sign of the wave vector
- section S12. Effect of antenna-to-antenna distance on nonreciprocity
- fig. S1. Nonreciprocal spin wave emission in MSSW mode.
- fig. S2. Nonreciprocity ratios in the time domain.
- fig. S3. FFT of spin wave signals.
- fig. S4. AMR measurements.
- fig. S5. VSM measurements.
- fig. S6. Out-of-plane hysteresis and nonreciprocity from simulations.
- fig. S7. Spatial and temporal magnetization excited by the antenna.
- fig. S8. Ta thickness dependence of spin wave resonance frequencies.
- fig. S9. Nonreciprocity in devices with an Ta overlayer.

fig. S10. Nonreciprocity from different directions of field sweep.

fig. S11. Nonreciprocity from opposite signs of the wave vector.

fig. S12. Spin wave packets in devices with a longer antenna-to-antenna distance (20 μm).

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