

## MATERIALS SCIENCE

Carbon-boron clathrates as a new class of  $sp^3$ -bonded framework materials

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Carbon-based frameworks composed of  $sp^3$  bonding represent a class of extremely lightweight strong materials, but only diamond and a handful of other compounds exist despite numerous predictions. Thus, there remains a large gap between the number of plausible structures predicted and those synthesized. We used a chemical design principle based on boron substitution to predict and synthesize a three-dimensional carbon-boron framework in a host/guest clathrate structure. The clathrate, with composition  $2Sr@B_6C_6$ , exhibits the cubic bipartite sodalite structure (type VII clathrate) composed of  $sp^3$ -bonded truncated octahedral  $C_{12}B_{12}$  host cages that trap  $Sr^{2+}$  guest cations. The clathrate not only maintains the robust nature of diamond-like  $sp^3$  bonding but also offers potential for a broad range of compounds with tunable properties through substitution of guest atoms within the cages.

## INTRODUCTION

As a fundamental building block of nature, carbon is unrivaled in its diversity to form stable structures with other elements and itself. One-dimensional (1D) carbon-based materials (e.g., polymers) have thoroughly reshaped society over the past century, in addition to providing the building blocks for life. 2D materials (e.g., graphene and metal borocarbides) attracted attention due to remarkable properties potentially useful for advanced technology (1–4). The limited number of 3D  $sp^3$  carbon-based structures includes diamond, lonsdaleite (a hexagonal diamond allotrope) (5), B-doped diamond (6), SiC (7), and  $BC_2N$  (8). These materials all have several attractive properties for applications that include hardness, strength, thermal conductivity, and electron mobility. Boron carbide also contains  $sp^3$ -hybridized carbon, but these atoms serve as dopants within or linkages between B icosahedra (9, 10), rather than establishing the overall structural framework. Of broader interest are 3D covalent organic frameworks (COFs), which are formed by linking  $sp^2$ -hybridized molecular building blocks, that have attracted attention for gas storage and separations (11). Compared with the exquisite synthetic control over porous COF materials, the experimental progress in denser  $sp^3$  carbon-based structures lags far behind.

Numerous 3D carbon allotropes and compounds are predicted to have feasible energies for synthesis and interesting properties

(12). However, it remains unclear whether any of these materials can be produced in the laboratory. Aside from the diamond structure, almost no other  $sp^3$  carbon-based frameworks are known or can be stabilized at atmospheric pressure. For example, longstanding predictions of 3D  $sp^3$ -bonded  $C_3N_4$  networks have not been realized thus far (13), and high-pressure polymeric phases of  $CO_2$ , which consist of a network of  $CO_4$  tetrahedra, decompose into molecular  $CO_2$  phases when decompressed (14).

Another impressive 3D  $sp^3$  material is the carbon clathrate. Carbon-based clathrates are open-framework structures composed of host cages that trap guest atoms, in which all host atoms are linked by four-coordinate bonds. As  $sp^3$ -bonded frameworks, carbon-based clathrates potentially represent strong and lightweight materials that also offer tunable properties through manipulation of the occupancy and type of guest atoms within the cages. Despite their prominence in other systems with tetrahedral coordination (15–18), carbon-based clathrates are tremendously challenging to synthesize. Attempts to synthesize carbon clathrates go back at least 50 years since they were postulated following the formation of inorganic silicon clathrates (18, 19), and their possible structures and properties are of long-standing interest (20, 21). Some proposed carbon clathrates are expected to exhibit exceptional mechanical properties with tensile and shear strengths exceeding those of diamond (22), while large electron-phonon coupling is predicted to give rise to conventional superconductivity with high transition temperatures (23, 24). If produced, these materials would represent a class of diamond-like compounds wherein the electronic structure is tunable by adjusting the occupancy of electron-donating (or withdrawing) atoms within the cages (25).

A substantial amount of research has been performed to answer the persisting question of whether carbon clathrate structures are accessible by experiment. First-principles density functional theory (DFT) calculations indicate that both filled and guest-free carbon clathrates are energetically unfavorable, but by energies as low as 0.07 eV/atom relative to diamond (for reference, commercially produced  $C_{60}$  is metastable by nearly six times that energy) (12, 26). Synthesis of carbon clathrates might therefore proceed through a nonequilibrium pathway (e.g., formation from a high-energy precursor or deposition method) or through a chemical substitution/doping strategy to modify the intrinsic thermodynamic stability.

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The closest synthesized materials that have a structure only resembling a carbon clathrate are 3D polymers of  $C_{60}$  (27).

While nonequilibrium synthesis pathways remain feasible in concept, another strategy is to substitute boron for carbon atoms within the cage frameworks of carbon clathrates. The electron-deficient nature of boron creates the ability to form complex chemical bonding with itself or carbon to stabilize polyhedra, such as the icosahedral units in molecular carborane clusters (28, 29). Zeng *et al.* (30) calculated that boron substitution improves the intrinsic thermodynamic stability of carbon clathrate frameworks. Nevertheless, no thermodynamically stable carbon-clathrate was predicted after examination of a small subset of possible B substitution schemes in Li-filled carbon clathrates. A broad search of potential B substitution schemes is needed to validate this chemical stabilization principle. We used automatic structure searching methods (31) to predict and then synthesize a thermodynamically stable carbon-boron  $sp^3$ -bonded clathrate.

## RESULTS

We conducted an extensive search in the Sr-B-C system (Fig. 1A) at pressures from 0 to 200 GPa after broader searching in other ternary B-C systems (see Materials and Methods for additional information). For the Sr-B-C system, we determined that several high-pressure compounds are thermodynamically stable with respect to elemental mixtures (Fig. 1B). At 50 GPa, we calculated that the hexagonal  $P6_3/mmc$  and  $\gamma$ -B structures are the most stable forms of Sr and B, respectively, while diamond is the most stable structure for C. We found the compounds  $Sr_5C_2$ , SrC,  $SrC_2$ ,  $Sr_5C_2$ , SrB,  $SrB_2$ ,  $SrB_4$ , and  $SrB_6$  are stable binaries on the convex hull at 50 GPa. We found no energetically stable B-C binary compounds above 50 GPa, in agreement with a previous computational study (32).

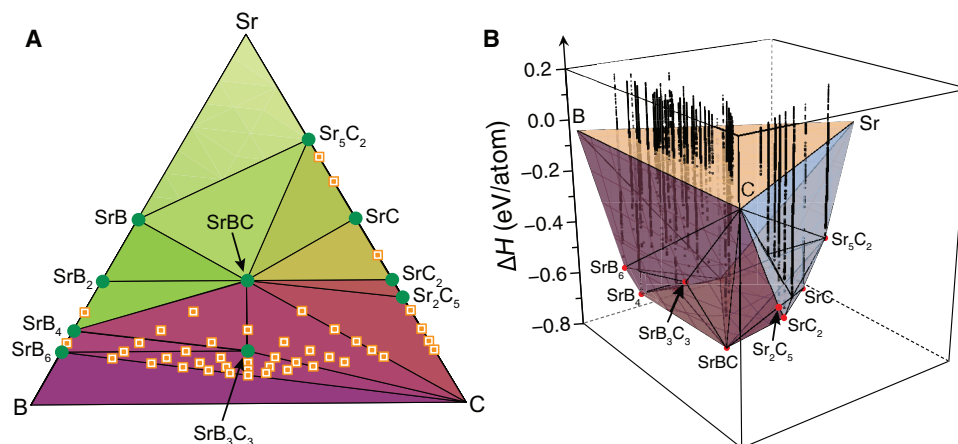
We predicted two stable ternary compounds at 50 GPa. The first hexagonal SrBC ( $P6_3/mmc$ ) exhibits 2D layers of six-membered B-C rings stacked between layers of  $Sr^{2+}$  cations, similar to intercalated graphite. This SrBC phase is isostructural with LiBC found at ambient pressure (3) and similar to other metal borocarbide phases with different ring tilings (2, 4). The second ternary compound is cubic ( $Pm\bar{3}n$ ) with composition  $2Sr@B_6C_6$  ( $SrB_3C_3$ ) and

has the type VII clathrate structure, known for the clathrate hydrate  $HPF_6 \cdot 6H_2O$  (33) and superhydrides such as  $CaYH_{12}$  (34). The topology of  $SrB_3C_3$  is that of bipartite sodalite (sod-b), which is distinguished from the sodalite structure (sod) in that carbon atoms are only bonded to boron atoms and vice versa.

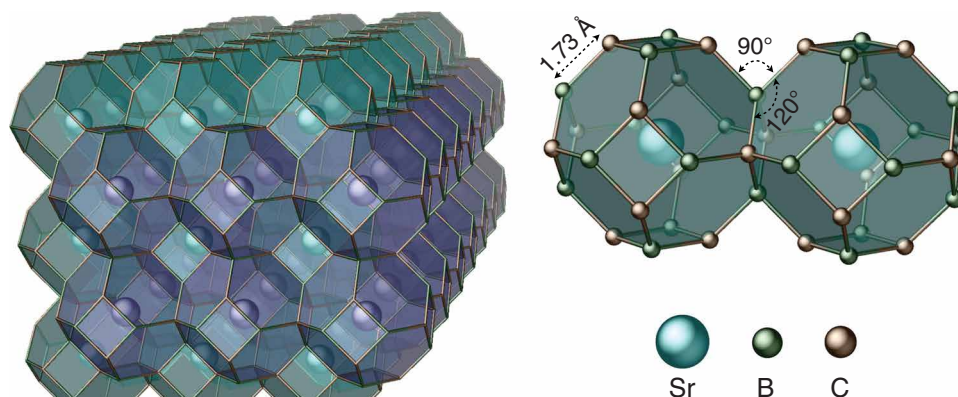
The  $SrB_3C_3$  clathrate framework (Fig. 2) is composed of a single truncated octahedral cage with six four-sided faces and eight six-sided faces ( $4^66^8$ ). The cages are composed of 24 vertices with alternating C and B atoms, and each cage contains a single  $Sr^{2+}$  cation at the center. We predicted that this type VII clathrate phase is thermodynamically stable from 50 to at least 200 GPa at 0 K. The material does not exhibit imaginary phonon frequencies at any pressure indicating dynamic stability and a favorable structure for metastable recovery to ambient conditions. At zero pressure, the optimized lattice parameter is 4.88 Å, and the structure contains one unique B-C bond length of 1.73 Å. The boron-doped clathrate is much more stable than its pure carbon counterpart. At 50 GPa,  $2Sr@C_{12}$  is metastable by 0.667 eV/atom, while  $2Sr@B_6C_6$  lies on the convex hull.

Our DFT calculations show that at 50 GPa,  $SrB_3C_3$  is a stable product of exothermic reactions of the pure elements and of readily accessible binary compounds. We therefore conducted diamond anvil cell (DAC) experiments using homogeneous fine-grained mixtures of  $SrB_6$ ,  $SrC_2$ , and glassy C targeting the stoichiometric reaction  $SrB_6 + SrC_2 + 4C \rightarrow 2SrB_3C_3$ . We conducted additional experiments using only mixtures of binary compounds, where the most energetically favorable reaction we calculated was for  $SrB_6 + 3SrC_2 \rightarrow SrB_3C_3 + 3SrBC$ . We compressed mixtures of the powders in Ne or  $Al_2O_3$  media and heated near ~2500 K using an infrared fiber laser, while synchrotron x-ray diffraction (XRD) patterns were collected to monitor structural changes in situ.

The starting  $SrB_6$  has the  $LaB_6$  ( $Pm\bar{3}n$ ) structure, whereas  $SrC_2$  takes on the acetylide structure of  $CaC_2$  ( $I4/mmm$ ). When compressed at room temperature,  $SrB_6$  remains in the starting cubic phase, and  $SrC_2$  transforms to the  $R\bar{3}m$  structure above 14 GPa, eventually appearing amorphous above 50 GPa. Upon heating near 2500 K, the intensities of diffraction peaks from the starting compounds vanish, and a series of new reflections appear. At 57(3) GPa, we indexed these sharp lines to a phase-pure BCC lattice with  $a = 4.5972(2)$  Å, in



**Fig. 1. Stable compounds in the Sr-B-C system.** (A) Ternary phase diagram at 50 GPa. Green circles represent thermodynamically stable compounds, while orange squares represent metastable compositions used in the search. (B) Ternary convex hull for the Sr-B-C system at 50 GPa based on formation enthalpies. Compounds with enthalpy data represented by red points are on the convex hull and thermodynamically stable against decomposition. Black points show the formation enthalpies of metastable structures found in the structure searches.



**Fig. 2. Structure of  $\text{SrB}_3\text{C}_3$  clathrate.** The cubic structure ( $Pm\bar{3}n$ ) is composed of face-sharing boron-carbon cages that encapsulate  $\text{Sr}^{2+}$  cations. Each cage contains 24 atoms with six four-sided faces and eight six-sided faces ( $4^66^8$ ). Different color cages are used to emphasize the stacking of cages that tile 3D space.

excellent correspondence with the type VII  $\text{SrB}_3\text{C}_3$  clathrate ( $Pm\bar{3}n$ ) with a calculated lattice parameter of  $a = 4.593 \text{ \AA}$  at the same pressure.

We compared the calculated XRD pattern of the  $\text{SrB}_3\text{C}_3$  clathrate with experimental scattering data (Fig. 3A). Given the nearly complete experimental powder averaging statistics, the quantitative diffraction intensities are representative of atomic positions and are in excellent agreement with the calculated pattern of  $\text{SrB}_3\text{C}_3$  clathrate to the experimental resolution limit of  $0.75 \text{ \AA}$ . All allowed reflections with appreciable intensity are observed to this limit, consistent with the formation of  $\text{SrB}_3\text{C}_3$  clathrate.

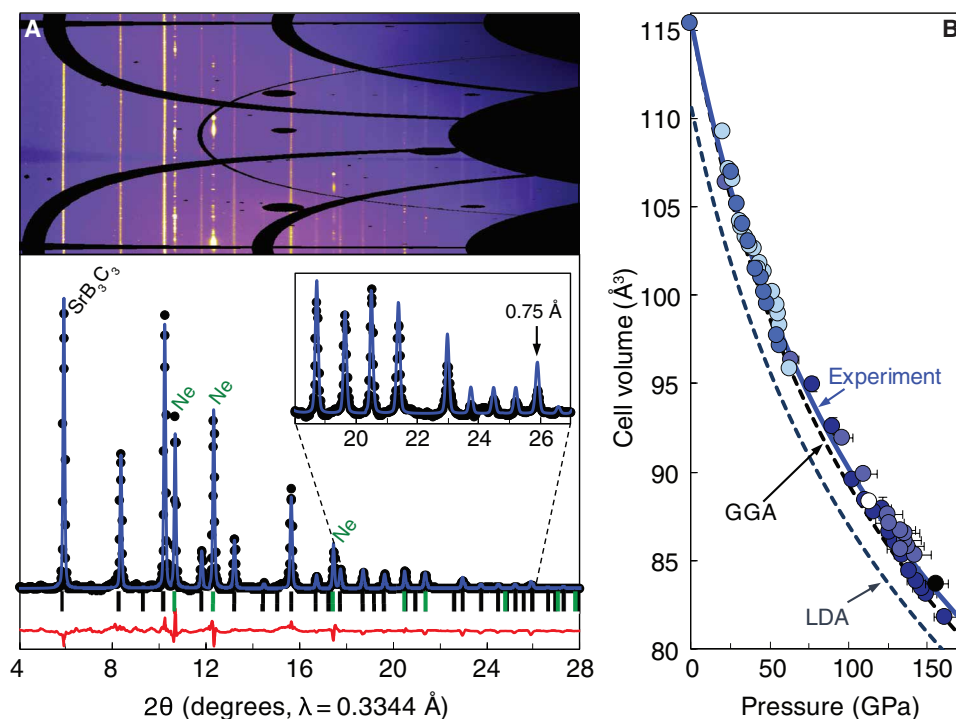
Given the large contrast in electron density between Sr and the framework atoms, the heavier element dominates the intensity of scattered x-rays. While the formation of  $\text{SrB}_3\text{C}_3$  is strongly supported by the stoichiometric conversion of the starting materials and agreement with the calculated stable structure, the intensities of the allowed reflections that differentiate the primitive bipartite structure from the BCC sodalite version are minimal, and it is challenging to confirm the clathrate structure on the basis of powder XRD alone. By annealing a sample at high-pressure and high-temperature conditions [ $56(4) \text{ GPa}$ ,  $\sim 3000 \text{ K}$ ], we were able to produce crystalline grains suitable for single-crystal diffraction. Single-crystal diffraction offers a more decisive characterization of the unit cell because it includes 3D geometrical relationships in addition to  $d$  spacings. Analysis of the data confirms the  $\text{sp}^3$ -bonded framework and allows us to unambiguously determine the clathrate crystal structure, exhibiting the symmetry of the  $Pm\bar{3}n$  space group. Structural models for the ordered bipartite structure are better than other models for the C/B framework (table S1). Furthermore, the formation of the clathrate phase is confirmed by the experimental  $P$ - $V$  equation of state, which uniquely distinguishes clathrate cage structures based on agreement with DFT calculations over a broad pressure range from 0 to  $150 \text{ GPa}$  (Fig. 3B).  $\text{SrB}_3\text{C}_3$  and other higher-energy cage variants are much less compressible than all other stable and metastable elemental structures and binary/ternary compounds calculated except pure boron allotropes and diamond (fig. S1). Energy-dispersive x-ray spectroscopy measurements performed on a recovered sample exhibiting the cubic XRD pattern are also consistent with the formation of  $\text{SrB}_3\text{C}_3$  clathrate near this composition (fig. S2).

The  $\text{SrB}_3\text{C}_3$  framework exhibits strong covalent bonding between  $\text{sp}^3$ -hybridized B and C atoms and weak interactions with the  $\text{Sr}^{2+}$  guest. This strong  $\text{sp}^3$ -hybridized covalent framework guarantees a

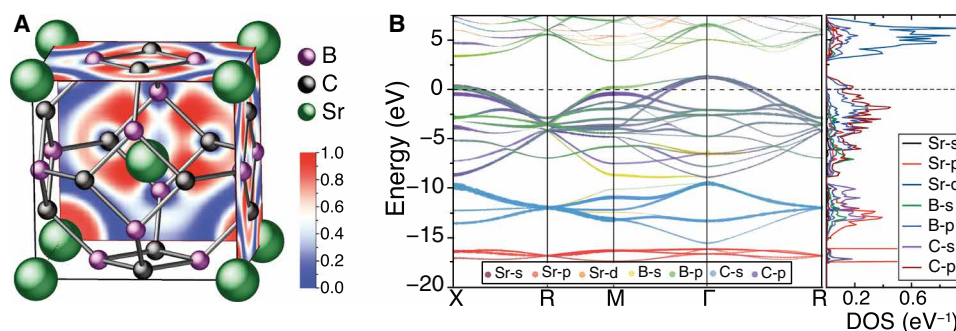
high value for the bulk modulus [experimental  $B_0 = 249(3) \text{ GPa}$ ] and incompressibility of  $\text{SrB}_3\text{C}_3$  clathrate. The estimated Vickers hardness for the  $\text{SrB}_3\text{C}_3$  clathrate, based on calculated elastic properties, is  $24 \text{ GPa}$ , comparable with the hardness of tungsten carbide (35). On the basis of electron count,  $\text{SrB}_3\text{C}_3$  should be a hole conductor, and calculations show that it is. All-carbon, four-coordinate zeolites are insulators at low pressures as closed-shell systems analogous to diamond. A sodalite all-carbon clathrate fits in this class, as does isoelectronic  $[\text{C}_3\text{B}_3]^{3-}$ .  $\text{SrB}_3\text{C}_3$  is one electron per formula unit short of this magic (insulator) electron count. The calculated band structure (Fig. 4) shows a good gap for one electron more than  $\text{SrB}_3\text{C}_3$ .

## DISCUSSION

$\text{SrB}_3\text{C}_3$  clathrate is likely the first member of a new class of strong and lightweight  $\text{sp}^3$ -bonded carbon-based frameworks with tunable properties. Because boron anions are isoelectronic to carbon atoms in the B-C framework, the bipartite structure should also exhibit exceptional properties due to the similar nature of bonding in hypothetical pure carbon cages.  $\text{SrB}_3\text{C}_3$  clathrate is calculated to be conducting, and samples exhibit a metallic luster. The electronic band structure (Fig. 4B) shows the coexistence of steep and flat bands close to the Fermi level (near the M point), which sets favorable conditions for superconductivity (36). Thus,  $\text{SrB}_3\text{C}_3$  is a candidate for phonon-mediated superconductivity at ambient pressure. Furthermore, the ability to trap various kinds of guest atoms in the  $\text{sp}^3$ -bonded cages allows the carbon-based clathrates to exhibit diamond-like mechanical properties with a tunable electronic structure. While  $\text{SrB}_3\text{C}_3$  is metallic, the electronic structure may be modified by substituting different guest atoms. For example,  $\text{LaB}_3\text{C}_3$ , which we also predict to be thermodynamically stable at high pressure, has a band gap due to the balanced electron count. The removal of guest ions from the cages offers an additional space to explore the possibility for guest-free structures (37). Although  $\text{SrB}_3\text{C}_3$  is only thermodynamically stable at high pressure, we recovered it at atmospheric pressure, similar to diamond. At  $1 \text{ atm}$ ,  $\text{SrB}_3\text{C}_3$  persists when kept under an inert atmosphere. The clathrate degrades when exposed to the moisture in air over a few hours. The clathrate we discovered provides a platform for investigating other substitutions of the framework or guest atoms. More broadly, using boron to stabilize carbon frameworks may be applicable to entirely different structure types.



**Fig. 3. XRD and equation of state of  $\text{SrB}_3\text{C}_3$ .** (A) Experimental XRD data (black points) collected at 57(3) GPa with Rietveld refinement (blue line) of the  $\text{SrB}_3\text{C}_3$  phase. Green ticks indicate contributions from Ne with Le Bail refinement. The 2D diffraction ("cake") aligned with the integrated pattern shows nearly complete powder averaging with sharp peaks for the  $\text{SrB}_3\text{C}_3$  phase. Black regions on the detector image indicate nonintegrated ("masked") regions due to diamond anvil reflections and features of the detector. The inset shows a magnified view at high angle with sharp  $\text{SrB}_3\text{C}_3$  peaks to a limiting resolution of 0.75 Å. (B) Experimental third-order Birch-Murnaghan equation of state (EoS) (solid blue line) with  $B_0 = 249(3)$  GPa,  $B_0' = 4.0$  (fixed) and calculated EoS (dashed lines) with  $B_0$  (DFT-LDA) = 257 GPa,  $B_0' = 4.0$  (fixed);  $B_0$  (DFT-GGA) = 225 GPa,  $B_0' = 4.0$  (fixed). Different colored symbols represent data points from six independent experimental runs.



**Fig. 4. Electronic properties of  $\text{SrB}_3\text{C}_3$  at 0 GPa.** (A) 2D electron localization function (ELF) for  $\text{SrB}_3\text{C}_3$ . The ELF indicates the probability of finding electrons in different regions of the crystal. Large ELF values (>0.6) indicate the formation of covalent bonds. (B) Electronic band structure for  $\text{SrB}_3\text{C}_3$  projected onto atomic orbitals represented by different colors, where the width of each band is proportional to the weight of the corresponding orbital character. The projected density of states (DOS) in  $\text{SrB}_3\text{C}_3$  is shown in the right. The Fermi energy is set to 0 eV (dashed line).

## MATERIALS AND METHODS

### Calculations

Structure-searching calculations were performed using the CALYPSO structure prediction method (31) based on the global minimization of free energy using ab initio total energy calculations. Structure searching simulations through CALYPSO code were performed in the Sr-B-C system (including pure elements, binaries, and ternaries from  $\text{SrB}_x\text{C}_y$ , with  $0 \leq x$  and  $y \leq 6$ ) from 0 to 200 GPa after broader searching in ternary B-C systems with a variety of metals including Li, Na, Mg, and Ca. Total energy calculations were performed in the framework of DFT within the Perdew-Burke-Ernzerhof (38) general-

ized gradient approximation as implemented in the VASP (Vienna Ab Initio Simulation Package) code (39). The projector-augmented wave (PAW) method (40) was adopted with the PAW potentials taken from the VASP library, where  $4s^2 4p^6 5s^2$ ,  $2s^2 2p^1$ , and  $2s^2 2p^2$  are treated as valence electrons for Sr, B, and C atoms, respectively. The use of a plane-wave kinetic energy cutoff of 520 eV and dense  $k$ -point sampling, adopted here, was shown to give excellent convergence of total energies. Electronic charges were calculated using a Bader charge analysis scheme using a  $600 \times 600 \times 600$  fast Fourier transform grid. Phonon dispersion calculations were performed to determine the dynamical stability of the predicted structures by using

the finite displacement approach, as implemented in the Phonopy code (41). The Vickers hardness ( $H_V$ ) was estimated to be 24 GPa by using the empirical method proposed by Guo *et al.* (42).

### Synthesis

Strontium metal was purified by sublimation (950°C, dynamic vacuum) onto Mo foil from a graphite crucible; graphite powder was pretreated at 950°C under dynamic vacuum for 16 hours to remove adsorbed species. Strontium carbide ( $\text{SrC}_2$ ) was prepared by heating 2.780 g of Sr (31.7 mmol) with 0.371 g of graphite powder (30.9 mmol) in a capped graphite crucible under dynamic vacuum at 825°C for 16 hours. Excess Sr was subsequently sublimed at 950°C, and the resulting pale gray  $\text{SrC}_2$  powder (1.400 g, 81% yield) was isolated from the crucible under Ar. Powder XRD showed a small SrO impurity.  $\text{SrB}_6$  (EPSI Metals, 99.5%) and glassy carbon (Sigma-Aldrich, 99.95%) were purchased commercially and used without further purification. Binary ( $\text{SrB}_6 + 3\text{SrC}_2$ ) and ternary ( $\text{SrB}_6 + \text{SrC}_2 + 4\text{C}$ ) mixtures were prepared under an inert Ar atmosphere, sealed, and then vigorously milled using  $\text{Si}_3\text{N}_4$  media at 600 rpm for 1-min cycles over ~12 hours. The milled powders were removed from the media in an Ar glovebox, and thin plates (10  $\mu\text{m}$ ) were pressed between two diamond anvils with 1-mm culets and then loaded into DAC sample chambers using 100- to 300- $\mu\text{m}$  culets and Re gaskets. Ne or alumina plates served as the pressure transmitting media and thermal insulation from the diamond anvils. After being compressed to the target pressure between 50 and 150 GPa, samples were heated to ~2500 K using the double-sided laser heating systems at High Pressure Collaborative Access Team (HPCAT) (43) or GeoSoilEnviro Center for Advanced Radiation Sources (GSECARS) (44). Temperatures near 3000 K were needed to produce crystalline grains with size comparable to the x-ray beam.

### XRD and crystal structure determination

In situ powder XRD patterns were collected at the Advanced Photon Source, Sector 16, HPCAT, using a monochromatic wavelength of 0.4066 Å and at Sector 13, GSECARS, using a monochromatic wavelength of 0.3344 Å. The x-ray beam was focused on the sample, and scattered x-rays were detected using a PILATUS 1M or MarCCD detector. The sample-to-detector distance and geometrical parameters were calibrated using  $\text{CeO}_2$  and  $\text{LaB}_6$  standards with the DIOPTAS software (45). Pressure was calibrated using the equations of state of Ne and/or  $\text{Al}_2\text{O}_3$  and cross-checked using the SrO equation of state and ruby fluorescence in some samples. Rietveld refinements of XRD patterns were conducted using PowderCell (46) and GSAS with EXPGUI (47).

Diffraction from crystalline grains suitable for analysis by multigrain crystallography methods were analyzed in the Fable package (48) with ImageD11 and GrainSpotter (49). Indexing and integration were performed using GSE\_ADA/RSV (50), and the structure was solved using direct methods using SHELXS and refined using SHELXL-2014 (51) invoked within the Olex2 suite (52). A full description of the single-crystal diffraction analysis is provided in the Supplementary Materials. CSD 1949948 contains supplementary crystallographic data for this paper. These data can be obtained free of charge from the joint Cambridge Crystallographic Data Centre's and FIZ Karlsruhe's service to view and retrieve structures via [www.ccdc.cam.ac.uk/structures/](http://www.ccdc.cam.ac.uk/structures/).

### Scanning electron microscopy with energy-dispersive x-ray spectroscopy

Scanning electron microscopy (SEM) with energy-dispersive x-ray spectroscopy (EDS) was performed on a recovered sample using a

JEOL 6500F field-emission scanning electron microscope equipped with an Oxford Instruments X-Max detector (80  $\text{mm}^2$ ). The laser-heated DAC sample was recovered from ~50 GPa after synchrotron XRD measurements, which verified the cubic powder diffraction pattern of  $\text{SrB}_3\text{C}_3$ . The Re gasket was placed onto an Al stub with carbon adhesive under the insert atmosphere of an Ar glovebox. The stub was sealed under Ar and transferred into the microscope by using a polyethylene glove bag that was sealed around the exchange chamber and purged with Ar. A beam of 15 keV and 1 nA was focused to a working distance of 10 mm, and an energy range of 10 keV was divided into 2048 channels for an energy resolution of ~5 eV. Aztec software developed by Oxford Instruments was used for EDS data analysis.

### SUPPLEMENTARY MATERIALS

Supplementary material for this article is available at <http://advances.sciencemag.org/cgi/content/full/6/2/eaay8361/DC1>

Single-crystal diffraction analysis

Fig. S1. Stability of  $\text{SrB}_3\text{C}_3$  and analysis of different possible stoichiometries.

Fig. S2. SEM and EDS measurements.

Fig. S3. Raw XRD patterns.

Fig. S4. Sr-B-C phase identification.

Fig. S5. Optical images of  $\text{SrB}_3\text{C}_3$  during synthesis near 50 GPa.

Fig. S6. Electronic structures for  $\text{SrB}_3\text{C}_3$ .

Fig. S7. Phonon dispersion curves and energetic stabilities as a function of pressure.

Fig. S8. Fourier difference map ( $F_{\text{obs}} - F_{\text{calc}}$ ) from single-crystal analysis.

Table S1. Comparison of single-crystal diffraction refinement quality indicators for different

"colorings" of clathrate framework in crystal structure models of the  $\text{SrB}_3\text{C}_3$  clathrate.

Table S2. Calculated Bader partial charges of the  $\text{SrB}_3\text{C}_3$  clathrate at 0 GPa.

Table S3. Calculated structural parameters of Sr-B-C phases.

Table S4.  $\text{SrB}_3\text{C}_3$  lattice parameters during decompression.

Data file S1. CIF file of the single-crystal XRD data and the best model of the crystal structure for the  $\text{SrB}_3\text{C}_3$  clathrate.

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