

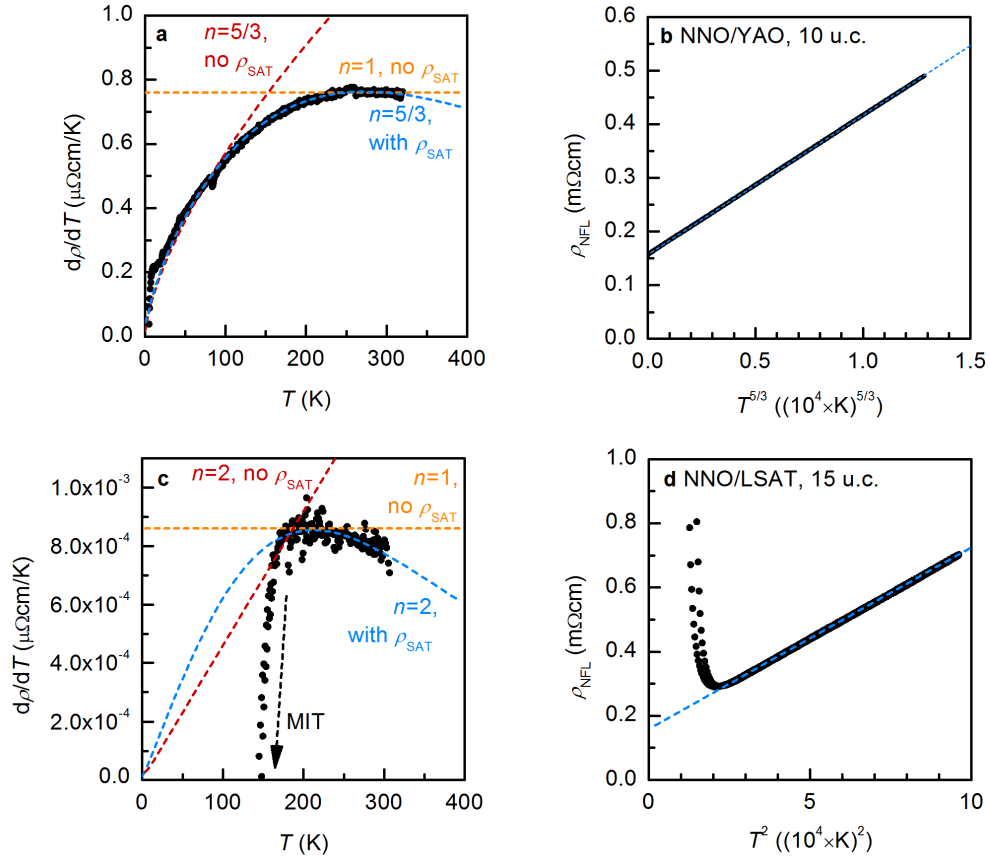
Supplementary Materials for  
**Tuning bad metal and non-Fermi liquid behavior in a Mott material:  
Rare-earth nickelate thin films**

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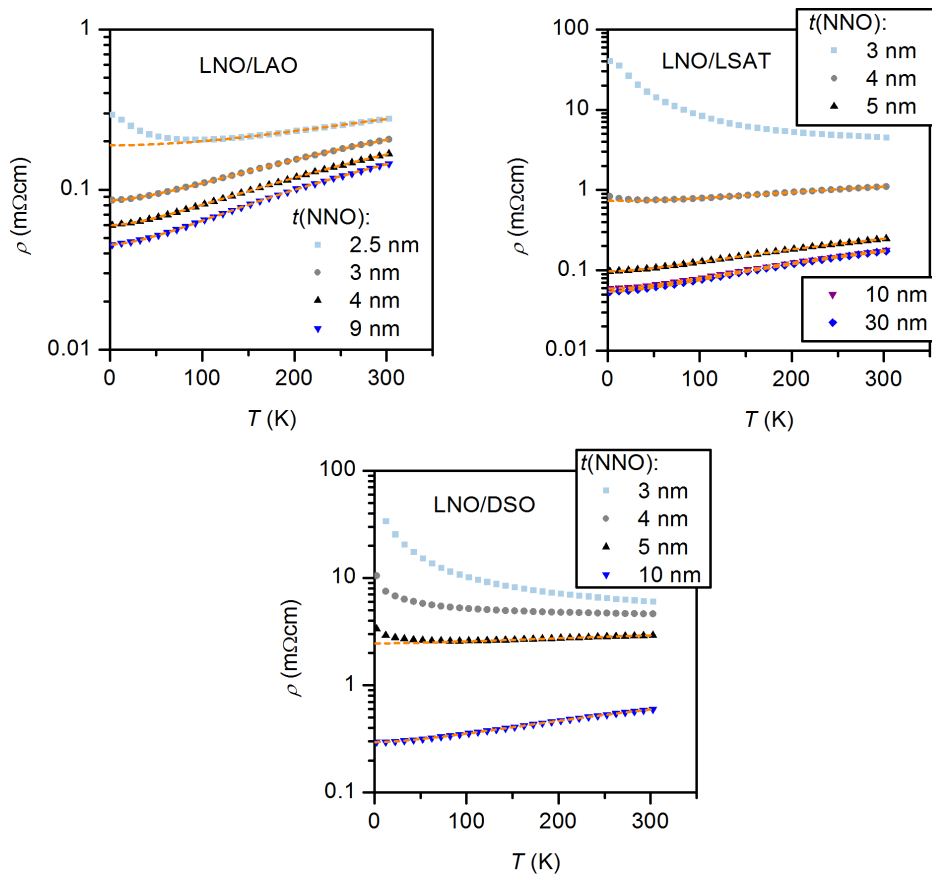
Published 6 November 2015, *Sci. Adv.* **1**, e1500797 (2015)  
DOI: 10.1126/sciadv.1500797

**The PDF file includes:**

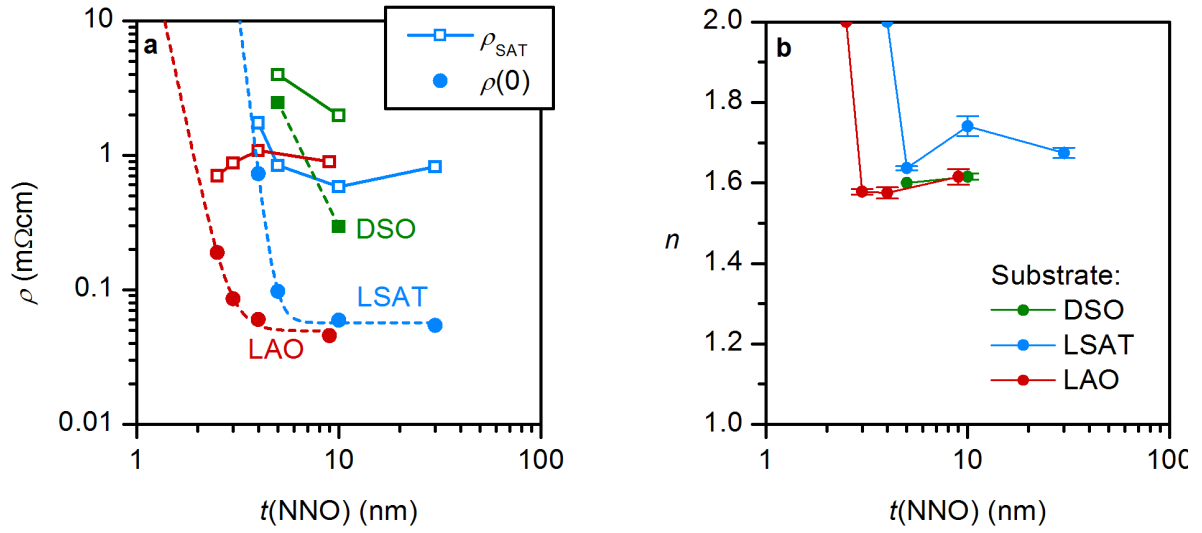
- Fig. S1. Resistivity as a function of temperature.
- Fig. S2.  $\rho$ - $T$  data for  $\text{LaNiO}_3$ .
- Fig. S3. Saturation resistivity and NFL behavior in  $\text{LaNiO}_3$ .
- Fig. S4. Electron-electron scattering coefficient  $A$ .



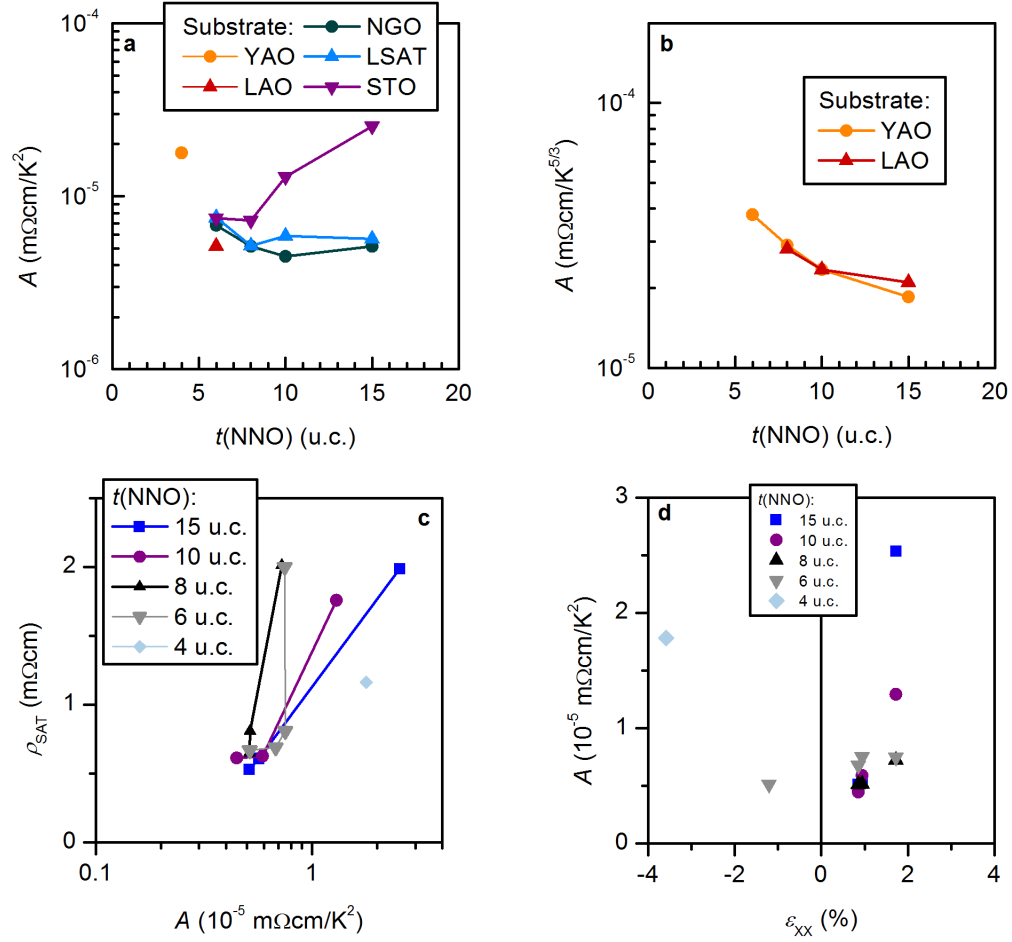
**Figure S1 | Resistivity as a function of temperature.** (a, c) Temperature derivative of the resistivity as a function of temperature for a 10 u.c. film on  $\text{YAlO}_3$ , a fully metallic non-Fermi liquid and a 15 u.c. thick film on LSAT, a LFL with a sharp metal-insulator transition near 150 K. The corresponding  $\rho_{\text{NFL}}$  as a function of  $T^n$  are shown in (b) and (d).



**Figure S2 |  $\rho$ - $T$  data for  $\text{LaNiO}_3$ .** Data re-plotted from ref. (33). The dashed orange lines are fits to Eqs. (1) and (2).



**Figure S3 | Saturation resistivity and NFL behavior in LaNiO<sub>3</sub>.** (a)  $\rho_{\text{SAT}}$  and  $\rho(0)$  extracted from data in Fig. S3. The condition  $\rho(0) = \rho_{\text{SAT}}$  accurately predicts the transition to an insulator at all temperatures. (b) Extracted exponent  $n$ , showing that thick LaNiO<sub>3</sub> films are non-Fermi liquids. All data are from ref. (33).



**Figure S4 | Electron-electron scattering coefficient  $A$ .** (a)  $A$  as a function of thickness for the LFLs ( $n = 2$ ). (b)  $A$  as a function of thickness for all NFL films ( $n = 5/3$ ). (c) Scaling between  $A$  and  $\rho_{\text{SAT}}$ , each curve corresponds to a specific  $\text{NdNiO}_3$  thickness,  $t_{\text{NNO}}$ . (d)  $A$  as a function of strain for different film thicknesses.